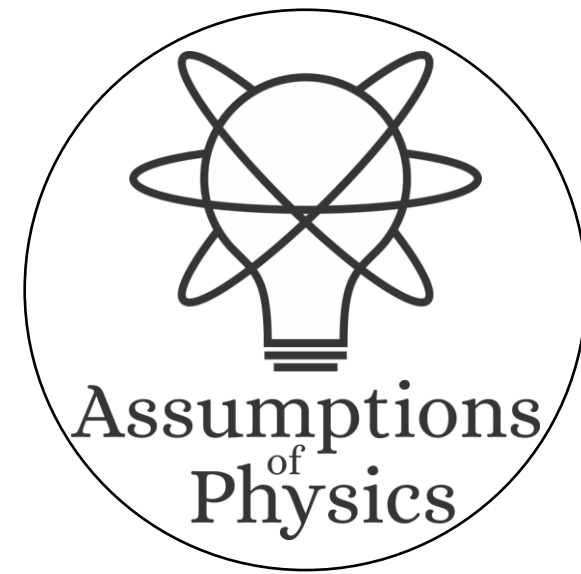


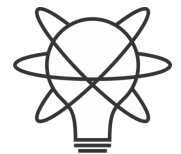
Assumptions of Physics
Summer School 2025
Reverse Physics
for classical mechanics

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Overview



<https://assumptionsofphysics.org/>

Assumptions
of
Physics

Assumption IR (Infinitesimal Reducibility). *The state of the system is reducible to the state of its infinitesimal parts. That is, specifying the state of the whole system is equivalent to specifying the state of its parts, which in turn is equivalent to specifying the state of its subparts and so on.*

Classical state is a distribution over the states of infinitesimal parts (i.e. particles)

Recovers pairs of conjugate variables for each DOF, frame-invariant entropy and count of states

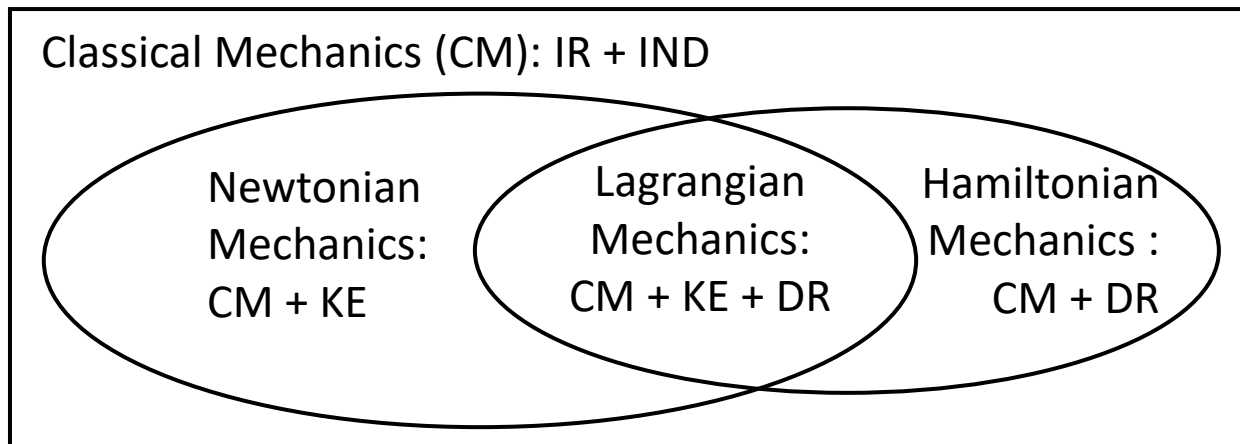
Assumption IND (Independent DOFs). *The system is decomposable into independent degrees of freedom. That is, the variables that describe the state can be divided into groups that have independent definition, units and count of states.*

Assumption DR (Determinism and Reversibility). *The system undergoes deterministic and reversible evolution. That is, specifying the state of the system at a particular time is equivalent to specifying the state at a future (determinism) or past (reversibility) time.*

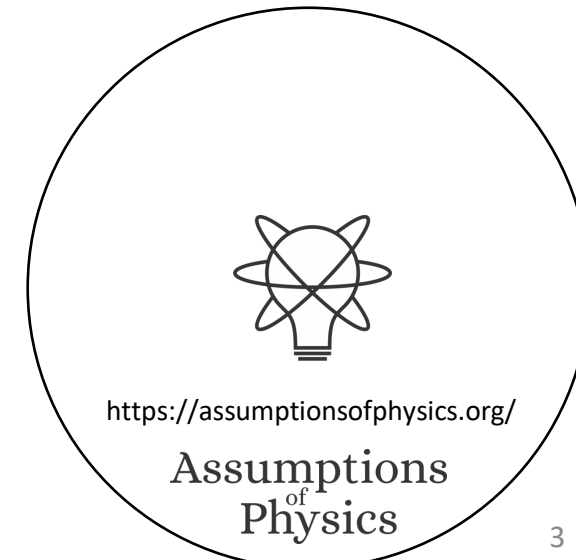
Conservation of entropy/count of states and DOF independence recovers Hamiltonian evolution

Links trajectories in space with dynamics, recovering inertia, masses and scalar/vector potential forces

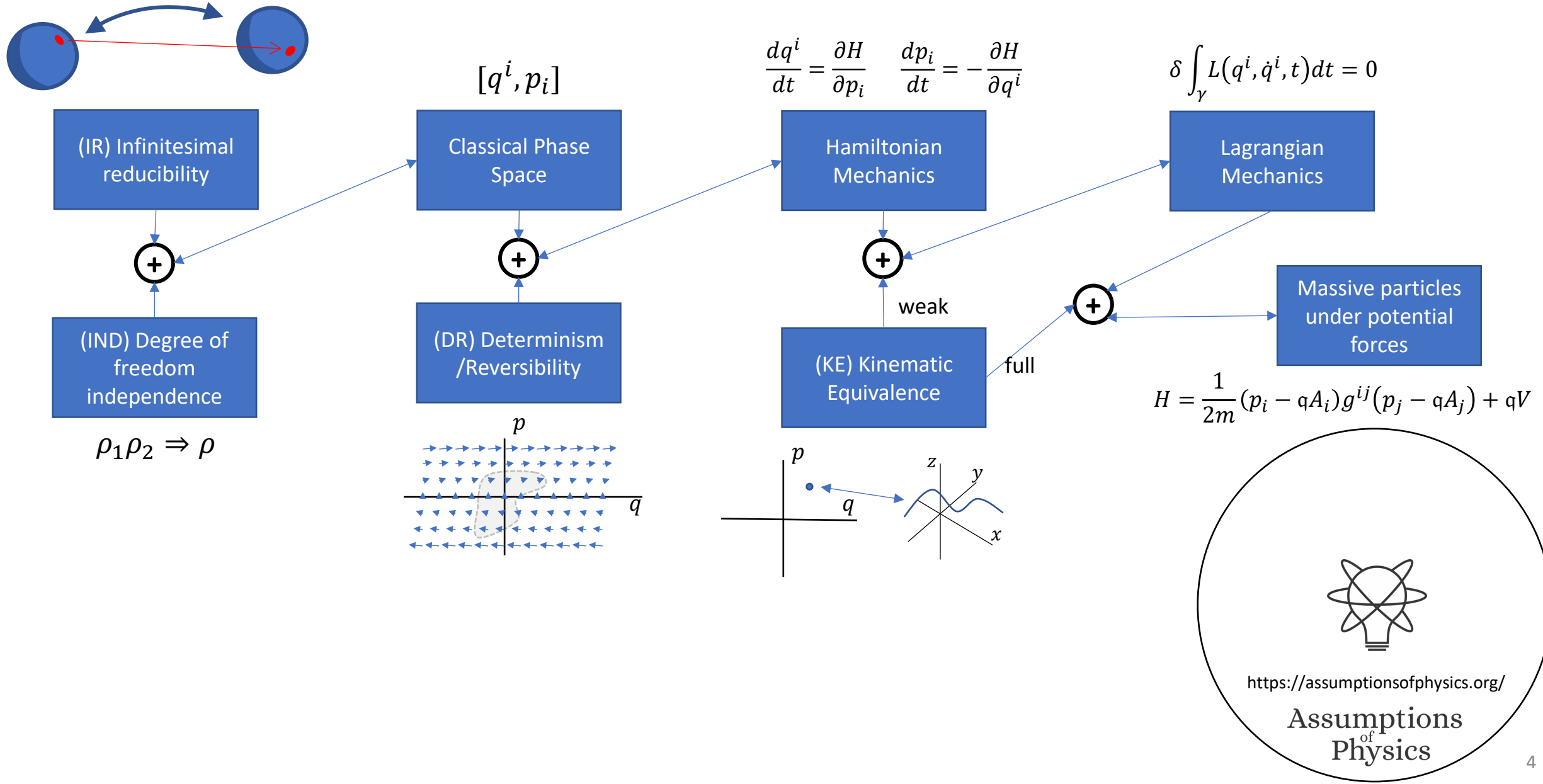
Assumption KE (Kinematic Equivalence). *The kinematics of the system is sufficient to reconstruct its dynamics and vice-versa. That is, specifying the motion of the system is equivalent to specifying its state and evolution.*



Assumptions for classical mechanics



Assumptions of classical mechanics



Each assumption can be expressed in multiple different but equivalent ways

Assumption DR (Determinism and Reversibility). *The system undergoes deterministic and reversible evolution. That is, specifying the state of the system at a particular time is equivalent to specifying the state at a future (determinism) or past (reversibility) time.*

The displacement field is divergenceless: $\partial_a S^a = 0$	(DR-DIV)
The Jacobian of time evolution is unitary: $ \partial_b \hat{\xi}^a = 1$	(DR-JAC)
Densities are conserved through the evolution: $\hat{\rho}(\hat{\xi}^a) = \rho(\xi^b)$	(DR-DEN)
Volumes are conserved through the evolution: $d\hat{\xi}^1 \dots d\hat{\xi}^n = d\xi^1 \dots d\xi^n$	(DR-VOL)
The evolution is deterministic and reversible.	(DR-EV)
The evolution is deterministic and thermodynamically reversible	(DR-THER)
The evolution conserves information entropy	(DR-INFO)
The evolution conserves the uncertainty of peaked distributions	(DR-UNC)

Assumption KE (Kinematic Equivalence). *The kinematics of the system is sufficient to reconstruct its dynamics and vice-versa. That is, specifying the motion of the system is equivalent to specifying its state and evolution.*

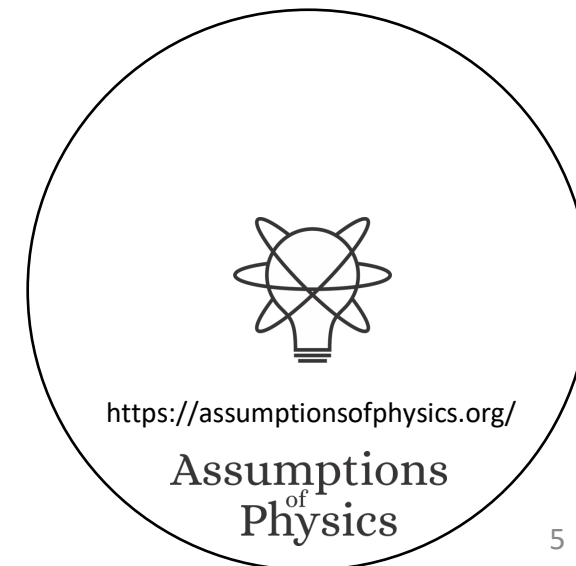
There is a linear relationship between conjugate momentum and velocity	(FKE-LIN)
The system under study is a massive particle under scalar and vector potential forces	(FKE-POT)
The Jacobian of the transformation between state variables and kinematic variables is a non-singular function of position only.	(FKE-NSIN)
Densities over phase space can be expressed in terms of position and velocity by rescaling the value at each point: $\rho(x^i, v^j) J(x^i) = \rho(q^i, p_j)$.	(FKE-DEN)
Areas and volumes in phase space can be expressed in kinematic variables, and the transformation depends on position only: $dx^1 \dots dx^n dv^1 \dots dv^n = J(x^i) dq^1 \dots dq^n dp_1 \dots dp_n$.	(FKE-VOL)
The symplectic form ω_{ab} can be expressed in kinematic variables, and its components are a linear function of velocity.	(FKE-SYMP)
Density expressed in velocity at the same position is proportional to the density over states	(FKE-PROP)
At each position, there exists a local inertial frame	(FKE-INER)
The position fully defines the units of all state variables, therefore an invertible transformation between momentum and velocity	(FKE-UNIT)
Uniform distributions along momentum correspond to uniform distributions along velocity	(FKE-UNIF)

Assumption IND (Independent DOFs). *The system is decomposable into independent degrees of freedom. That is, the variables that describe the state can be divided into groups that have independent definition, units and count of states.*

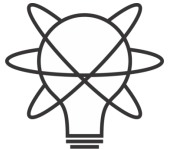
The system is decomposable into independent DOFs	(IND-DOF)
The system allows statistically independent distributions over each DOF	(IND-STAT)
The system allows informationally independent distributions over each DOF	(IND-INFO)
The system allows peaked distributions where the uncertainty is the product of the uncertainty on each DOF	(IND-UNC)

Assumption IR (Infinitesimal Reducibility). *The state of the system is reducible to the state of its infinitesimal parts. That is, specifying the state of the whole system is equivalent to specifying the state of its parts, which in turn is equivalent to specifying the state of its subparts and so on.*

The state of a classical system is given by a distribution over phase space.	(IR-DIST)
A classical system can be thought of as being made of infinitesimal parts, called particles.	(IR-INF)



Action principle one DOF



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Assumptions
of
Physics

Required background

- Hamiltonian mechanics

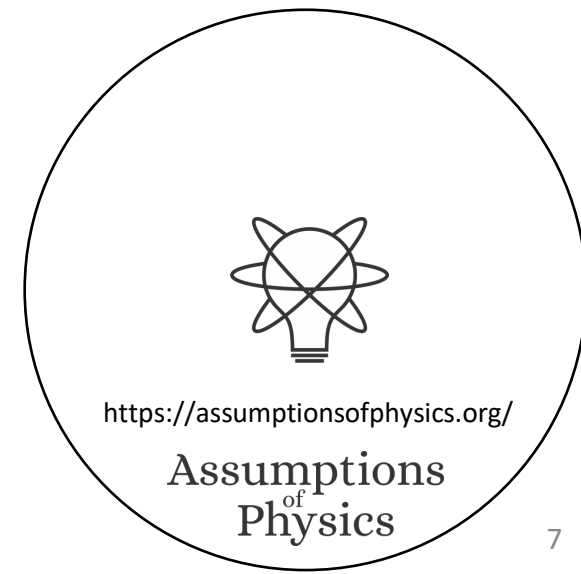
- Hamilton's equations: $\frac{dq}{dt} = \frac{\partial H}{\partial p} \quad \frac{dp}{dt} = -\frac{\partial H}{\partial q}$

- Lagrangian mechanics

- Action is the integral of the Lagrangian over a given path: $\mathcal{A}[\gamma] = \int_{\gamma} L(q(t), \dot{q}(t), t) dt$
 - Actual evolutions are the paths that make the action stationary: $\delta \mathcal{A} = 0$
 - Lagrangian is the Legendre transform of the Hamiltonian: $L = p\dot{q} - H$

- Vector calculus

- Given a vector field $\vec{v}$, a field line is a line always tangent to the field
 - The divergence $\vec{\nabla} \cdot \vec{v}$ is the flow of the field through an infinitesimal closed region
 - A divergence-free field $\vec{\nabla} \cdot \vec{B} = 0$ admits a vector potential $\vec{A}$ such that $\vec{B} = \vec{\nabla} \times \vec{A}$



$$\text{Action: } \mathcal{A}[\gamma] = \int_{\gamma} L(q(t), \dot{q}(t), t) dt$$

$$\text{Lagrangian: } L = T - U \quad \text{Does not always work!}$$

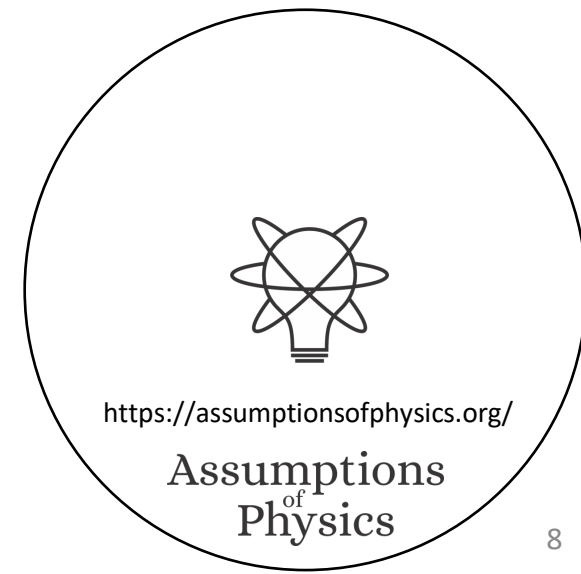
Lagrangian for particle with charge q :

$$L = \frac{1}{2} m v^2 + q \vec{v} \cdot \vec{A} - q V$$


 Kinetic


 ???


 Potential



Displacement
field

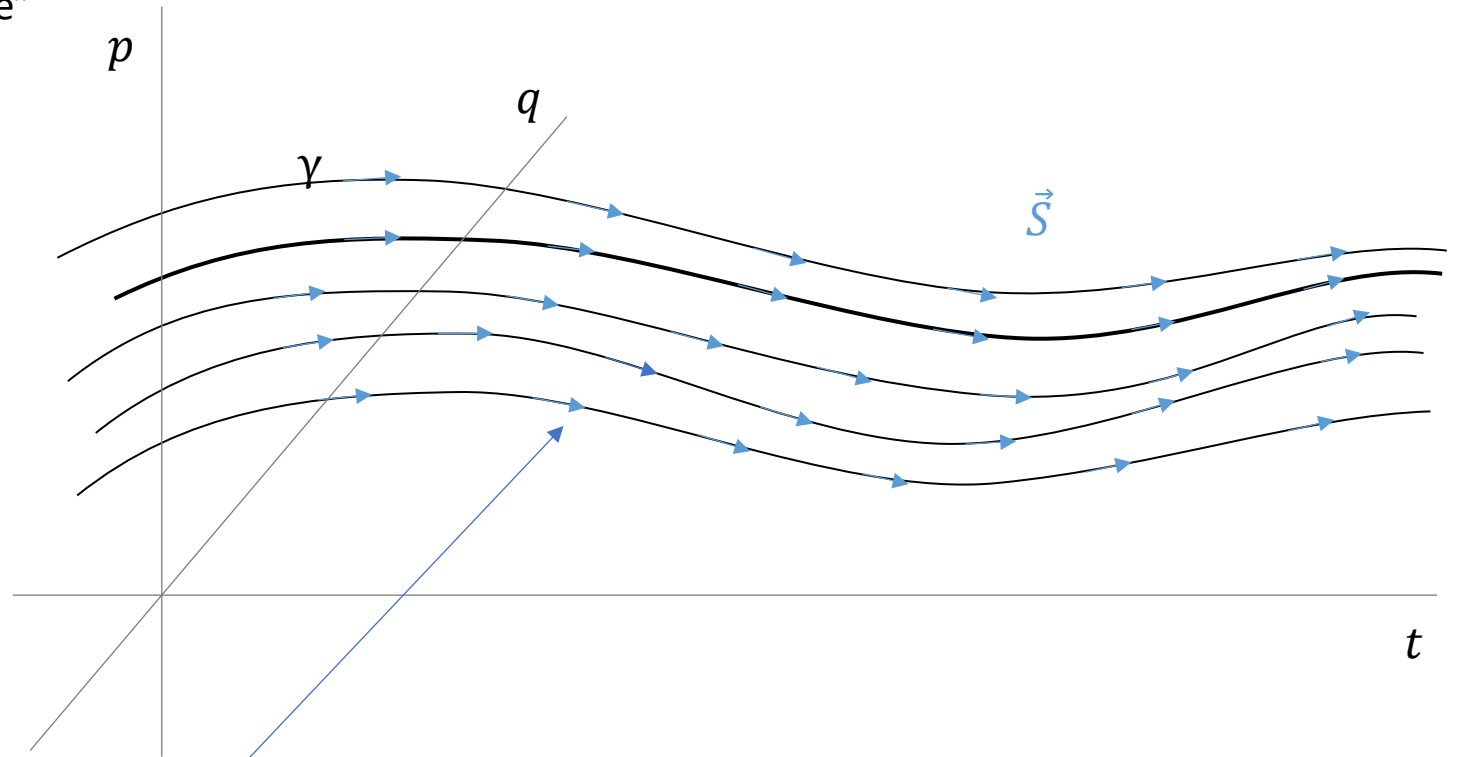
Tells us how states “move”

$$\vec{S} = \left[\frac{dq}{dt}, \frac{dp}{dt}, \frac{dt}{dt} \right]$$

Evolution

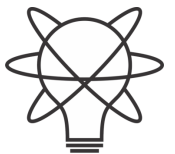
$$\gamma = [q(t), p(t), t]$$

Time is both a parameter
and a variable



Evolution is field lines
of the displacement field

Always tangent
to the field



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Assumptions
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(DR) Determinism and Reversibility

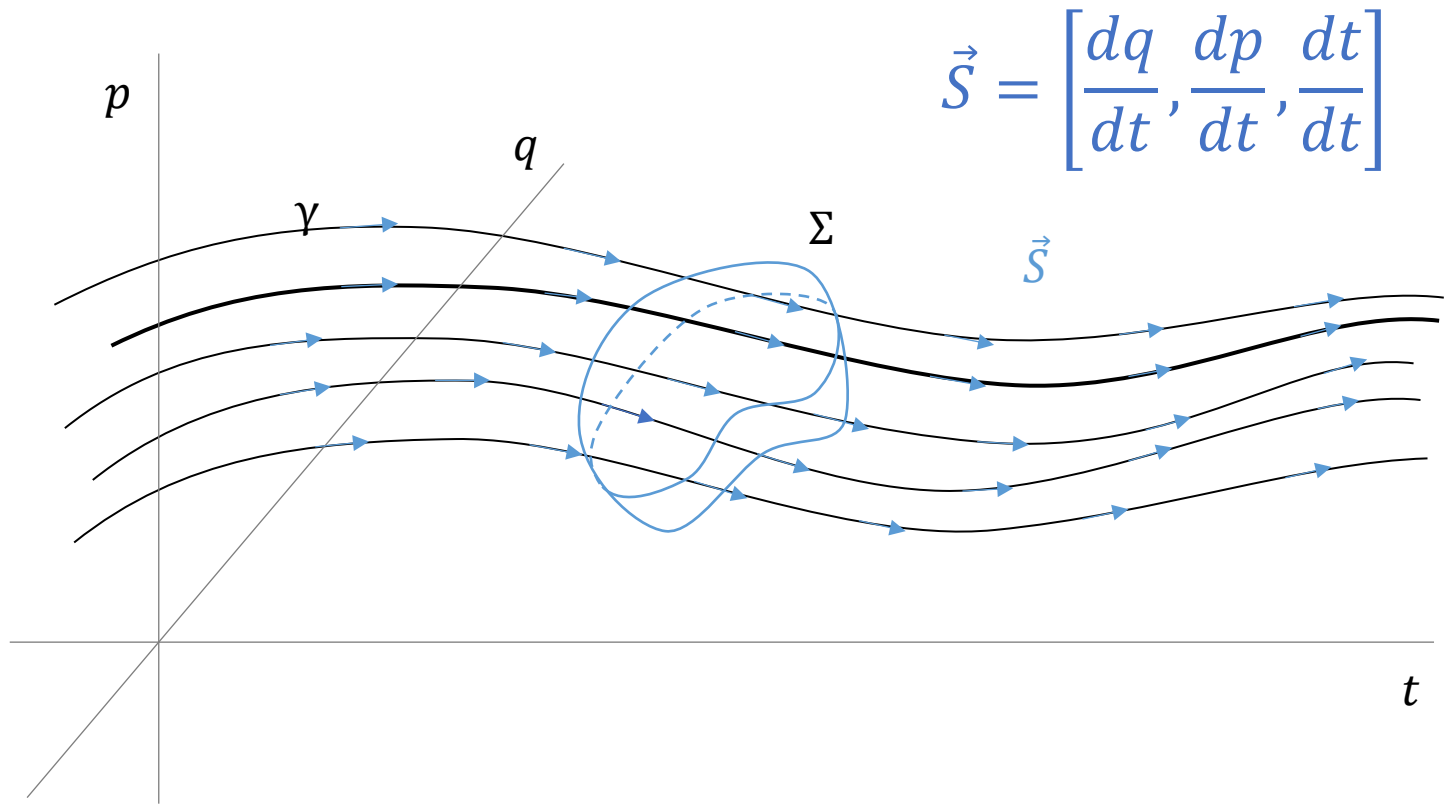
No states are lost or created
(count of states is preserved in time)

$$\oiint_{\Sigma} \vec{S} \cdot d\vec{\Sigma} = 0$$

⇒ Divergence-free displacement

$$\vec{\nabla} \cdot \vec{S} = 0$$

Displacement admits
a vector potential

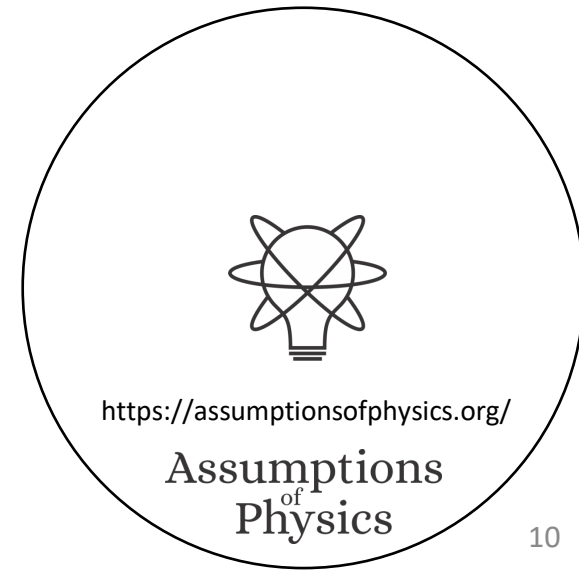


Vector potential

$$\vec{S} = -\vec{\nabla} \times \vec{\theta}$$

Minus sign added to match conventions

$$\vec{S} = \left[\frac{dq}{dt}, \frac{dp}{dt}, \frac{dt}{dt} \right]$$



$$\vec{\theta} = [\theta_q, \theta_p, \theta_t]$$

Fix gauge: $\theta_a \rightarrow \theta_a - \partial_a \int_0^p \theta_p dp$

$$\vec{\theta} = [\theta_q, 0, \theta_t]$$

$$S_t = \partial_p \theta_q - \partial_q \theta_p = \frac{dt}{dt} = 1$$

$$\vec{\theta} = [p, 0, \theta_t]$$

Rename θ_t

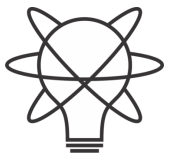
$$\vec{\theta} = [p, 0, -H]$$

Without loss of generality

$$\vec{S} = -\vec{\nabla} \times \vec{\theta}$$

$$\vec{\theta} = [p, 0, -H]$$

similar to four-momentum



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Assumptions
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$$\vec{S} = -\vec{\nabla} \times \vec{\theta}$$

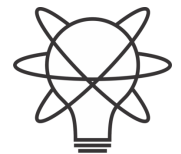
$$\vec{\theta} = [p, 0, -H(q, p)]$$

$$S^q = \frac{dq}{dt} = -\left(\frac{\partial}{\partial p}\theta_t - \frac{\partial}{\partial t}\theta_p\right) = \frac{\partial H}{\partial p}$$

Hamilton's equations

$$S^p = \frac{dp}{dt} = -\left(\frac{\partial}{\partial t}\theta_q - \frac{\partial}{\partial q}\theta_t\right) = -\frac{\partial H}{\partial q}$$

$$S^t = \frac{dt}{dt} = -\left(\frac{\partial}{\partial q}\theta_p - \frac{\partial}{\partial p}\theta_q\right) = 1$$

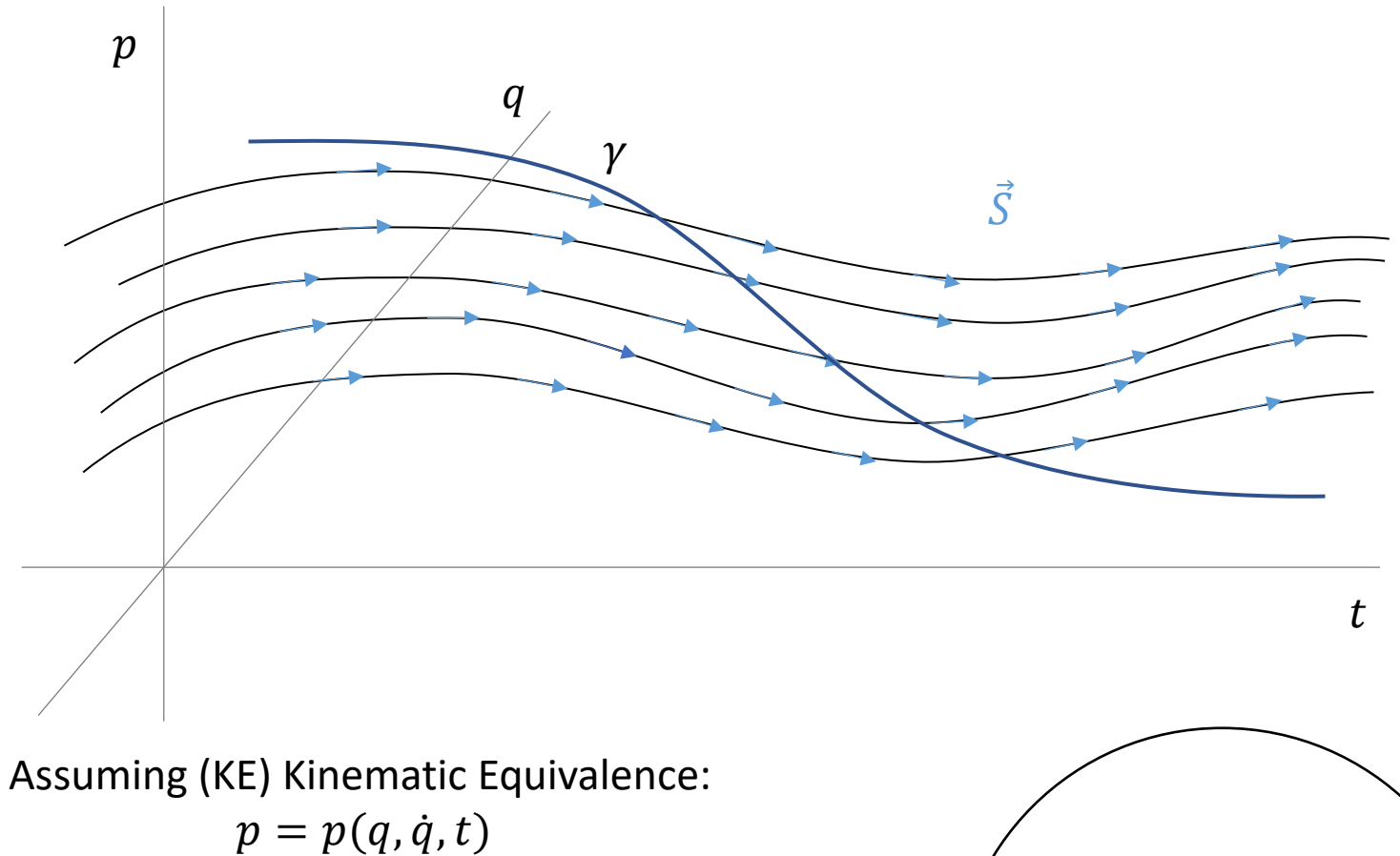


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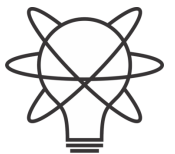
Assumptions
of
Physics

Line integral of $\vec{\theta}$

$$\begin{aligned} & \int_{\gamma} \vec{\theta} \cdot d\vec{\gamma} \quad \text{Express in components} \\ &= \int_{\gamma} [p, 0, -H] \\ & \quad \cdot \left[\frac{dq}{dt}, \frac{dp}{dt}, \frac{dt}{dt} \right] dt \\ &= \int_{\gamma} \left(p \frac{dq}{dt} - H \right) dt \\ &= \int_{\gamma} L(q, \dot{q}, t) dt \\ &= \mathcal{A}[\gamma] \end{aligned}$$



The action is the line integral
of the vector potential of the flow of states

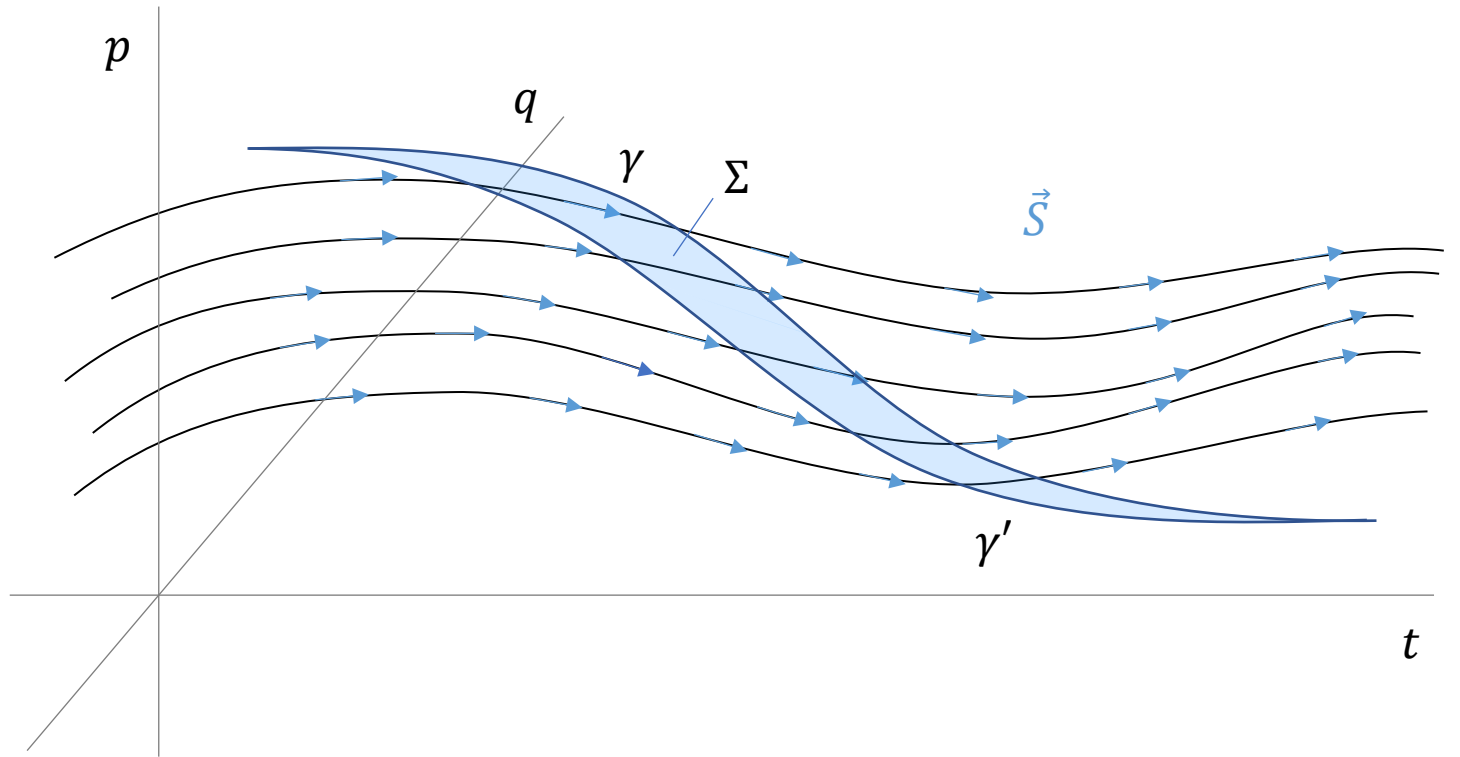


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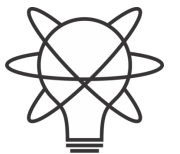
Assumptions
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Physics

Flow between a path γ
and its variation γ'

$$\begin{aligned}
 \Phi[\Sigma] &= - \iint_{\Sigma} \vec{S} \cdot d\vec{\Sigma} \\
 &= \iint_{\Sigma} \vec{\nabla} \times \vec{\theta} \cdot d\vec{\Sigma} && \text{Hamilton equations} \\
 &= \oint_{\partial\Sigma} \vec{\theta} \cdot d\vec{\gamma} && \text{Stokes' theorem} \\
 &= \int_{\gamma} \vec{\theta} \cdot d\vec{\gamma} - \int_{\gamma'} \vec{\theta} \cdot d\vec{\gamma} \\
 &= \delta \int_{\gamma} \vec{\theta} \cdot d\vec{\gamma} \\
 &= \delta \int_{\gamma} L dt \\
 &= \delta \mathcal{A}[\gamma]
 \end{aligned}$$



The variation of the action
is the flow between
path and its variation



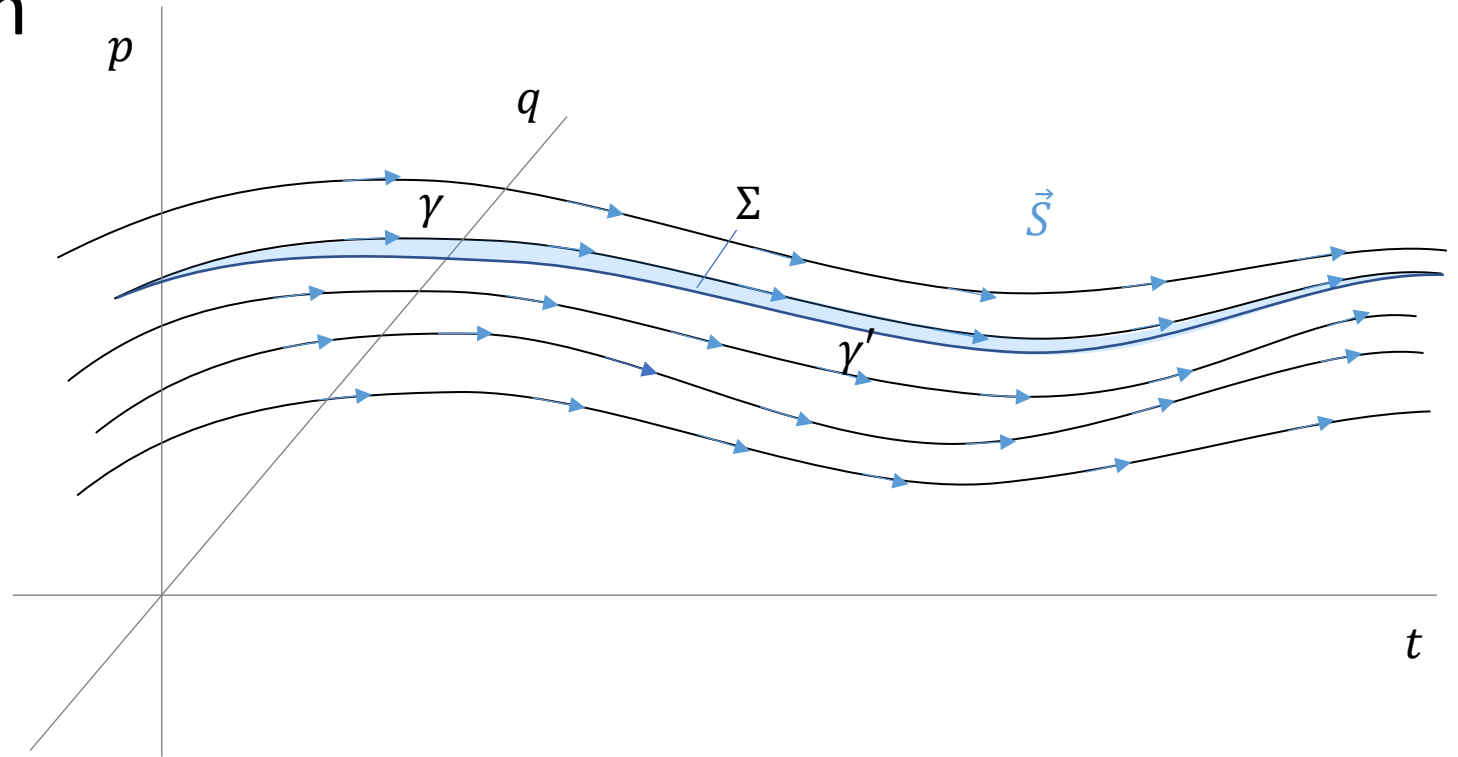
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Assumptions
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Physics

If the path γ is an evolution of the system

$$\begin{aligned}\Phi[\Sigma] &= - \iint_{\Sigma} \vec{S} \cdot d\vec{\Sigma} \\ &= \delta \int_{\gamma} \vec{\theta} \cdot d\vec{\gamma} = 0\end{aligned}$$

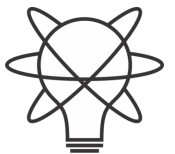
flow is tangent



If the path γ is not an evolution of the system, we can always find a variation γ' such that

$$- \iint_{\Sigma} \vec{S} \cdot d\vec{\Sigma} = \delta \int_{\gamma} \vec{\theta} \cdot d\vec{\gamma} \neq 0$$

Action is stationary over
and only over evolutions



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Assumptions
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Physics

1 DOF of classical phase space

+

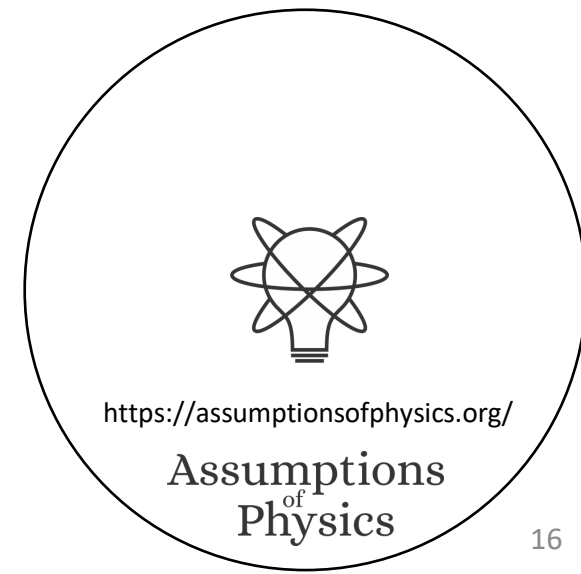
Determinism/Reversibility

+

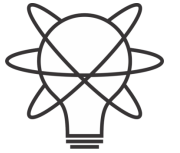
Kinematic equivalence

↕

Principle of stationary action



Action principle multiple DOF



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Assumptions
of
Physics

Required background

- Elements of differential geometry

- A differential n-form takes n vectors and returns a number that depends on the volume of the parallelepiped (must be anti-symmetric):

$$W[\gamma] = \int_{\gamma} dW = \int_{\gamma} F_a dx^a$$

$$\Phi[\Sigma] = \int_{\Sigma} d\Phi = \int_{\Sigma} B_{ab} v^a w^b$$

$$m[V] = \int_V dm = \int_V \rho_{abc} v^a w^b u^c$$

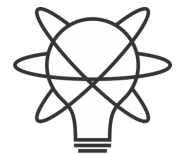
- Pseudo-vectors are really two-forms in 3D: $\vec{B} \rightarrow B_{ij}$
- The exterior derivative generalizes gradient, curl, divergence:

$$\nabla f \rightarrow \partial_a \wedge f \quad \nabla \times \vec{\theta} \rightarrow \partial_a \wedge \theta_b \quad \nabla \cdot \vec{B} \rightarrow \partial_a \wedge B_{bc}$$

- Generalized properties

- Generalized Stokes theorem: $\int_{\Sigma} \partial_a \wedge \theta_b d\Sigma^{ab} = \int_{\partial\Sigma} \theta_b d\gamma^a$
- Define closed form as having zero exterior derivative: $\partial_a \wedge \omega_{ab} = 0$
- A closed form (locally) admits a vector potential:

$$\omega_{ab} = \partial_a \wedge \theta_b = \partial_a \theta_b - \partial_b \theta_a$$



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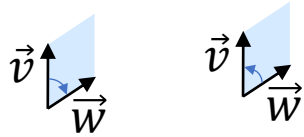
Assumptions
of
Physics

Characterize configuration flow

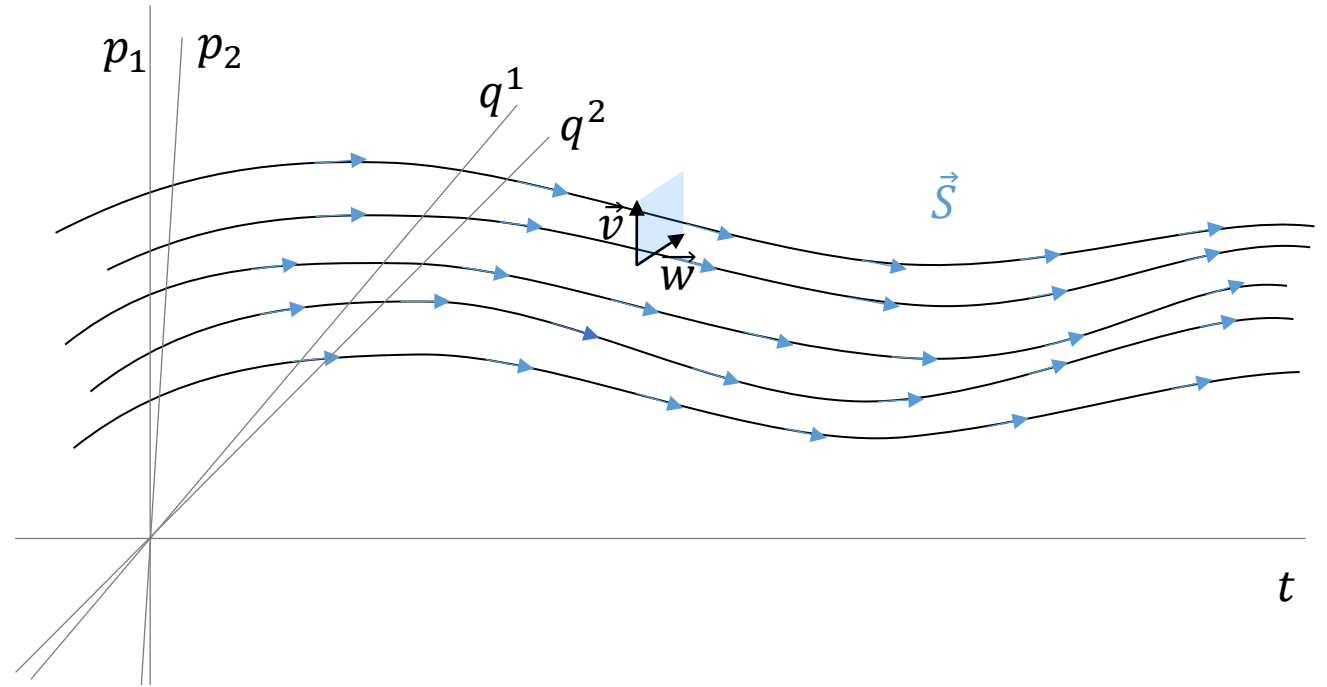
Flow of configurations through an infinitesimal surface

$$d\Phi = v^a \omega_{ab} w^b$$

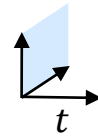
Function of surface: $\omega_{ab} = -\omega_{ba}$



change of order = change of orientation

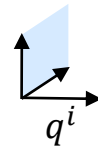


(DR) Determinism and reversibility: $\partial_t \wedge \omega_{ab} = 0$



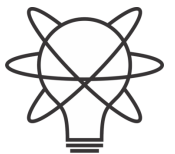
Configurations are conserved through time

(IND) DOF independence:
 $\partial_{q^i} \wedge \omega_{ab} = 0$
 $\partial_{p_i} \wedge \omega_{ab} = 0$



Configurations are conserved across DOF

IND + DR: ω_{ab} closed two-form ($\partial_a \wedge \omega_{bc} = 0$)



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Assumptions
of
Physics

$$\partial_a \wedge \omega_{bc} = 0$$

Closed form (DR) + (IND)

$$\omega_{ab} = -\partial_a \wedge \theta_b = -(\partial_a \theta_b - \partial_b \theta_a)$$

Vector potential

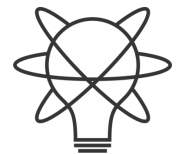
θ is the potential of ω

$$\begin{aligned} \partial_a \theta_b - \partial_b \theta_a &= \begin{bmatrix} \partial_{q^i} \theta_{q^j} - \partial_{q^j} \theta_{q^i} & \partial_{q^i} \theta_{p_j} - \partial_{p_j} \theta_{q^i} & \partial_{q^i} \theta_t - \partial_t \theta_{q^i} \\ \partial_{p_i} \theta_{q^j} - \partial_{q^j} \theta_{p_i} & \partial_{p_i} \theta_{p_j} - \partial_{p_j} \theta_{p_i} & \partial_{p_i} \theta_t - \partial_t \theta_{p_i} \\ \partial_t \theta_{q^j} - \partial_{q^j} \theta_t & \partial_t \theta_{p_j} - \partial_{p_j} \theta_t & \partial_t \theta_t - \partial_t \theta_t \end{bmatrix} \\ &= \begin{bmatrix} \partial_{q^i} p_j - \partial_{q^j} p_i & \partial_{q^i} 0 - \partial_{p_j} p_i & \partial_{q^i} (-H) - \partial_t p_i \\ \partial_{p_i} p_j - \partial_{q^j} 0 & \partial_{p_i} 0 - \partial_{p_j} 0 & \partial_{p_i} (-H) - \partial_t 0 \\ \partial_t p_j - \partial_{q^j} (-H) & \partial_t 0 - \partial_{p_j} (-H) & \partial_t (-H) - \partial_t (-H) \end{bmatrix} \\ &= \begin{bmatrix} 0 - 0 & 0 - \delta_i^j & -\partial_{q^i} H - 0 \\ \delta_j^i - 0 & 0 - 0 & -\partial_{p_i} H - 0 \\ 0 + \partial_{q^j} H & 0 + \partial_{p_j} H & -\partial_t H + \partial_t H \end{bmatrix} \\ &= \begin{bmatrix} 0 & -\delta_i^j & -\partial_{q^i} H \\ \delta_j^i & 0 & -\partial_{p_i} H \\ \partial_{q^j} H & \partial_{p_j} H & 0 \end{bmatrix}. \end{aligned}$$

$$\theta_a = [p_i \quad 0 \quad -H]$$

without loss of generality

$$\omega_{ab} = \begin{bmatrix} \omega_{q^i q^j} & \omega_{q^i p_j} & \omega_{q^i t} \\ \omega_{p_i q^j} & \omega_{p_i p_j} & \omega_{p_i t} \\ \omega_{t q^j} & \omega_{t p_j} & \omega_{t t} \end{bmatrix} = \begin{bmatrix} 0 & \delta_j^i & \partial_{q^i} H \\ -\delta_i^j & 0 & \partial_{p_i} H \\ -\partial_{q^j} H & -\partial_{p_j} H & 0 \end{bmatrix}$$



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Assumptions
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Physics

Displacement kills the flow

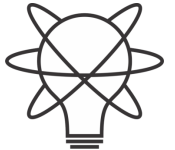
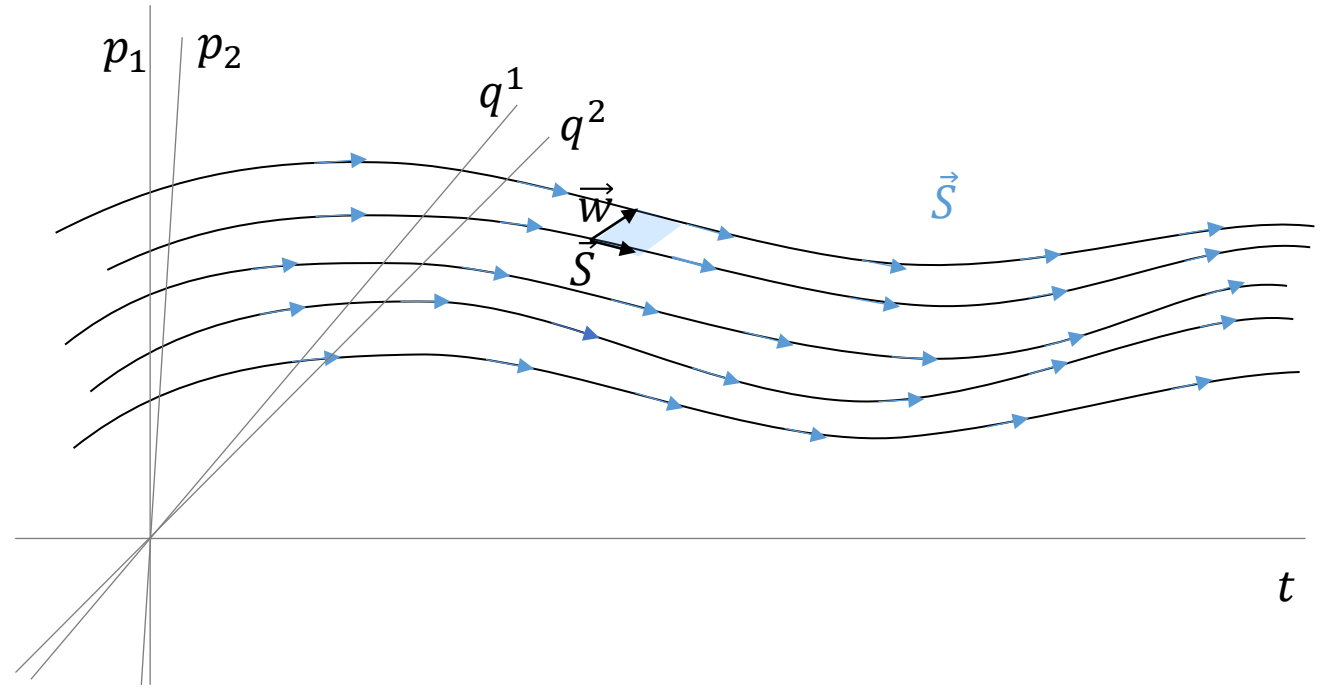
Displacement only direction that “kills” the flow

$$S^a \omega_{ab} w^b = 0 \quad \text{For all } \vec{w}$$

Flow is tangent (and only tangent) to the displacement

$$S^a \omega_{ab} = 0$$

Let's calculate!



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Assumptions
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Displacement kills the flow

$$S^a \omega_{ab} = 0$$

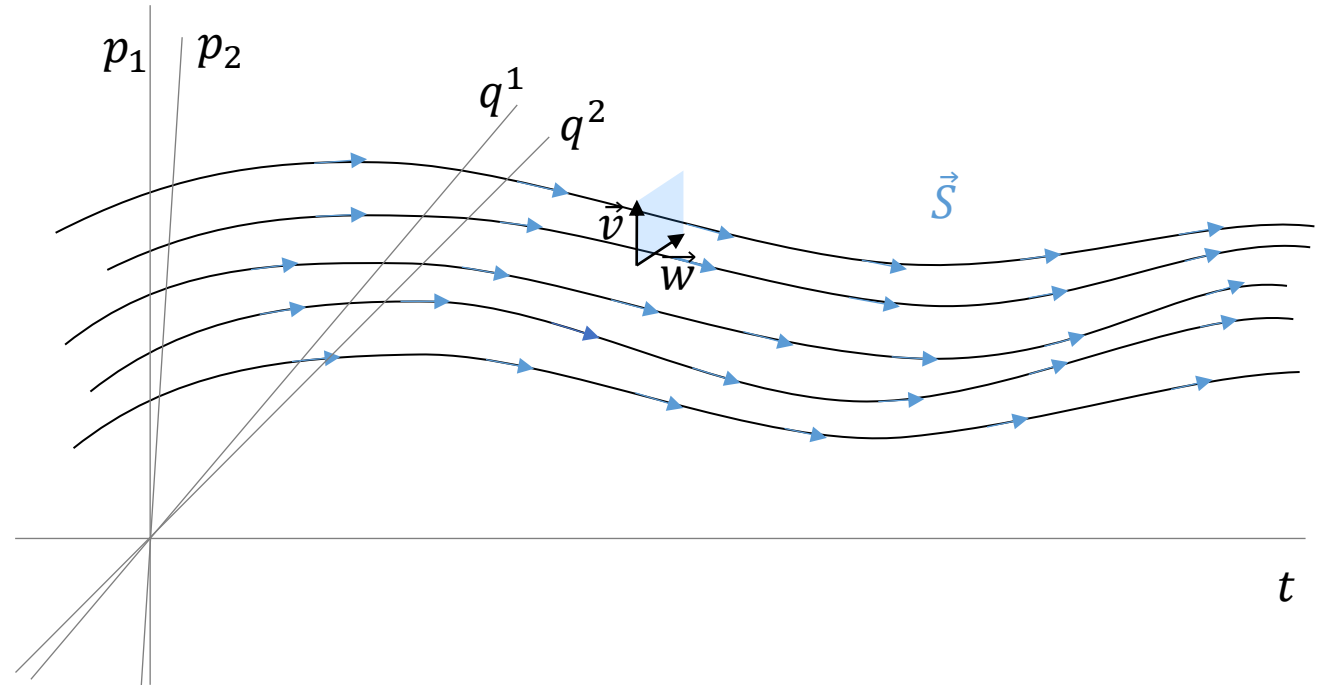
$$\omega_{ab} = \begin{bmatrix} 0 & \delta_j^i & \partial_{q^i} H \\ -\delta_i^j & 0 & \partial_{p_i} H \\ -\partial_{q^j} H & -\partial_{p_j} H & 0 \end{bmatrix}$$

$$\begin{aligned} S^a \omega_{aq^j} &= S^{q^i} \omega_{q^i q^j} + S^{p_i} \omega_{p_i q^j} + S^t \omega_{t q^j} \\ &= -S^{p_j} - S^t \partial_{q^j} H = -S^{p_j} - \partial_{q^j} H = 0 \end{aligned}$$

$$\begin{aligned} S^a \omega_{ap_j} &= S^{q^i} \omega_{q^i p_j} + S^{p_i} \omega_{p_i p_j} + S^t \omega_{t p_j} \\ &= S^{q^j} - S^t \partial_{p_j} H = S^{q^j} - \partial_{p_j} H = 0 \end{aligned}$$

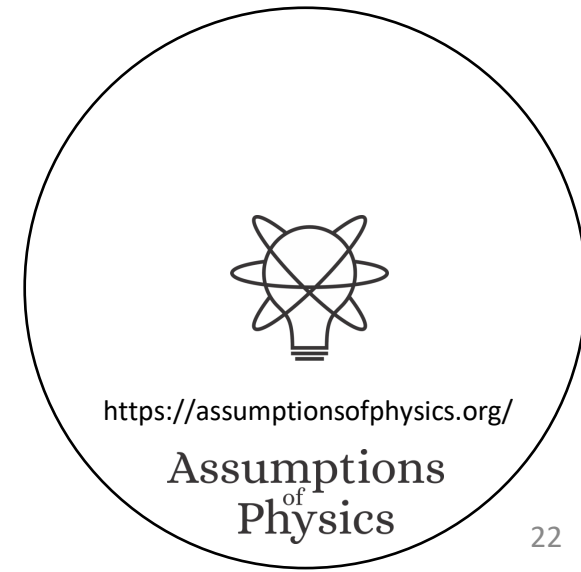
$$\begin{aligned} S^a \omega_{at} &= S^{q^i} \omega_{q^i t} + S^{p_i} \omega_{p_i t} + S^t \omega_{tt} \\ &= S^{q^i} \partial_{q^i} H + S^{p_i} \partial_{p_i} H \\ &= \partial_{p_i} H \partial_{q^i} H - \partial_{q^i} H \partial_{p_i} H = 0. \end{aligned}$$

Recovers Hamilton's equations



$$S^{p_i} = \frac{dp_i}{dt} = -\frac{\partial H}{\partial q^i}$$

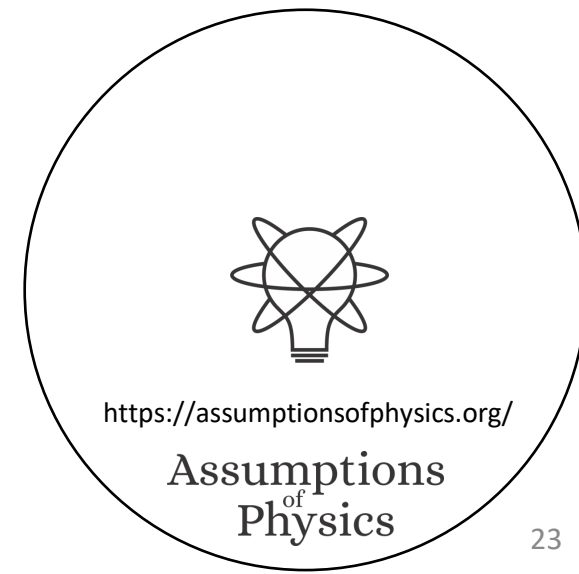
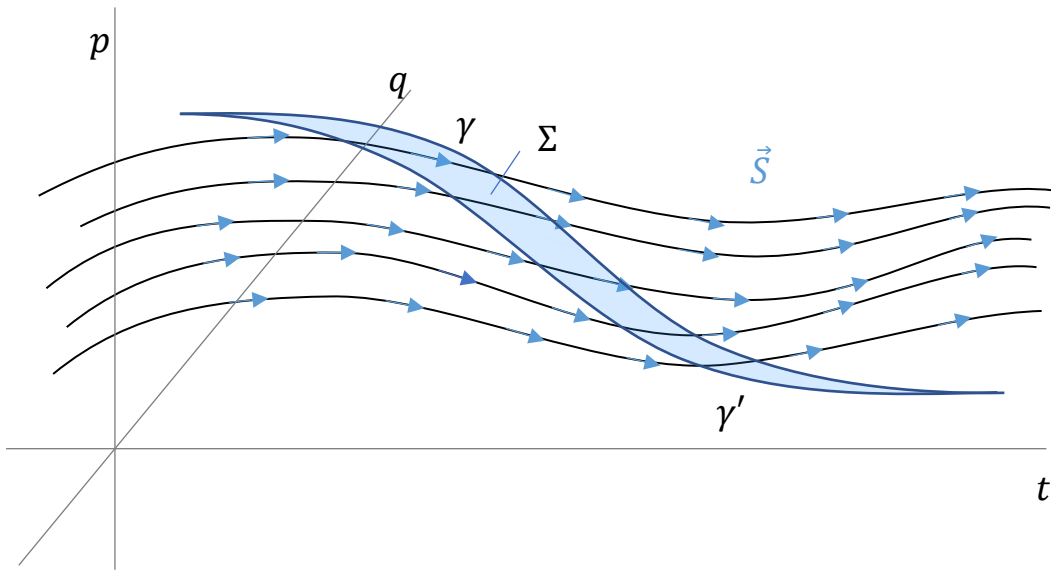
$$S^{q^i} = \frac{dq^i}{dt} = \frac{\partial H}{\partial p_i}$$



$$\xi^a = [q^i, p_i, t] \qquad \theta_a = [p_i, 0, -H]$$

$$\int_{\gamma} \theta_a d\gamma^a = \int_{\gamma} \theta_a d_t \xi^a dt = \int_{\gamma} (p_i d_t q^i - H) dt = \int_{\gamma} L dt = \mathcal{A}[\gamma]$$

Action is the line integral of the vector potential



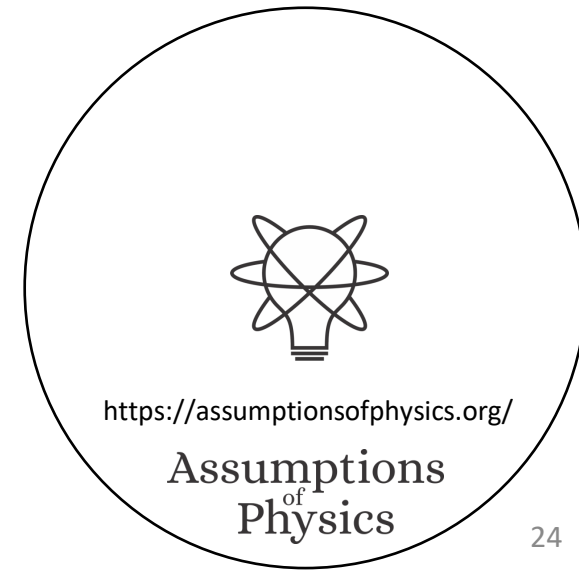
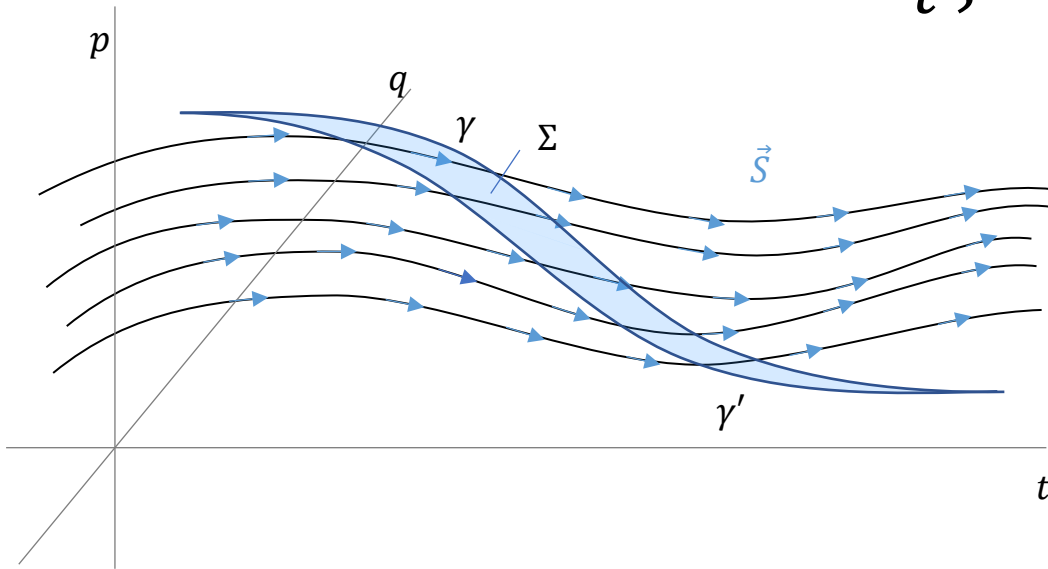
$$\delta \mathcal{A}[\gamma] = \delta \int_{\gamma} L dt = \delta \int_{\gamma} \theta_a d_t \xi^a dt = \oint_{\partial \Sigma} \theta_a d\xi^a = \int_{\Sigma} \partial_a \wedge \theta_b d\xi^a d\eta^b = - \int_{\Sigma} d\xi^a \omega_{ab} d\eta^b = -\Phi[\Sigma]$$

The variation of the action is the flow between
path and its variation

$$\delta \mathcal{A}[\gamma] = 0 \iff d\xi^a \omega_{ab} d\eta^b = dt d_t \xi^a \omega_{ab} d\eta^b = 0 \quad \forall d\eta^a$$

$$\iff d_t \xi^a \omega_{ab} = 0 \quad d_t \xi^a = S^a$$

Action is stationary
when the path kills the flow



DOF independence

+

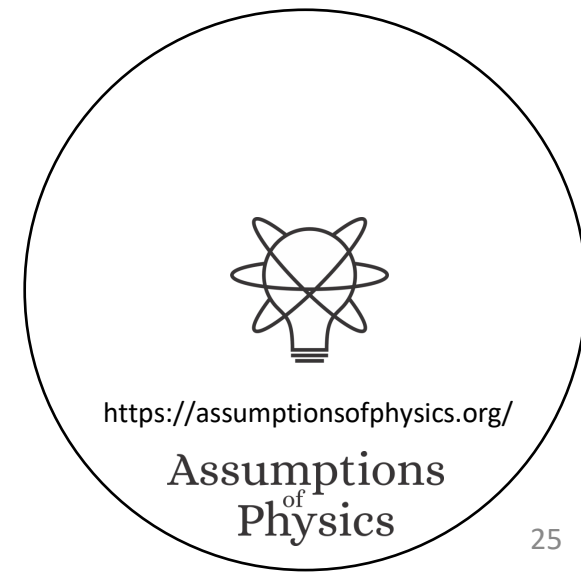
Determinism/Reversibility

+

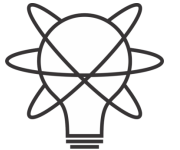
Kinematic equivalence

↕

Principle of stationary action



Relativistic mechanics



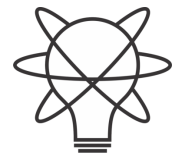
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Assumptions
of
Physics

Required background

- Special relativity

- Space and time coordinates are grouped together: $q^\alpha = [ct, q^i]$
- Metric tensor $g_{\alpha\beta}$ defines the geometry of space-time (i.e. lengths and angles)
- Using $(-, +, +, +)$ signature: $ds^2 = -dt^2 + dx^2 + dy^2 + dz^2$
- Energy and momentum are grouped together (four-momentum): $p_\alpha = \left[-\frac{E}{c}, p_i\right]$
- Four-velocity is the velocity with respect to proper time (time in the rest frame) and its norm is $-c^2$: $u^\alpha = d_\tau q^\alpha \quad u^\alpha g_{\alpha\beta} u^\beta = -c^2$
- For a free particle, relationship between four-momentum and four-velocity: $p_\beta g^{\beta\alpha} = p^\alpha = mu^\alpha$



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Assumptions
of
Physics

Want to extend Hamiltonian mechanics to handle change of coordinates
that mix space and time variables

$$q^\alpha = [t, q^i] \quad [q^i(t)] \rightarrow [t(s), q^i(s)]$$

Separate time-the-variable from time-the-parameter

Hamiltonian mechanics
on the extended
phase space

Recall

$$\theta_a = [p_i, 0, -H]$$

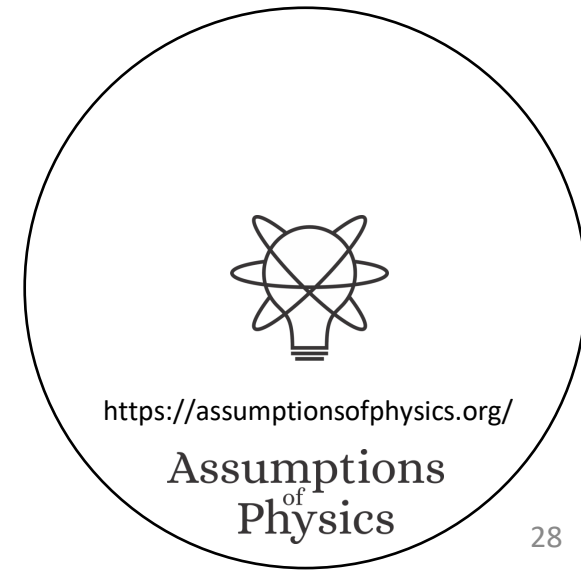
Form a covector

$$p_\alpha = [-E, p_i]$$

$\Rightarrow$ only possible choice for conjugates

$$\frac{dq^\alpha}{ds} = \frac{\partial \mathcal{H}}{\partial p_\alpha} \quad \frac{dp_\alpha}{ds} = -\frac{\partial \mathcal{H}}{\partial q^\alpha}$$

Hamiltonian constraint
 $\mathcal{H}(t, q^i, E, p_i) = 0$
Not just conserved quantity



$$q^\alpha = [t, q^i]$$

$$p_\alpha = [-E, p_i]$$

$$\omega = dq^i dp_i - dt dE$$

Energy-momentum covector and negative sign for temporal DOF appear in the extended phase space without knowledge of space-time geometry (or c)

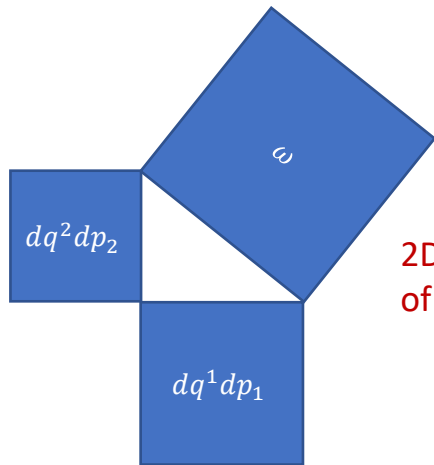
Lorentzian relativity is the only “correct” one

It extends the count of configurations per DOF to the temporal DOF in the correct way

$$\omega = dq^1 dp_1 + dq^2 dp_2$$

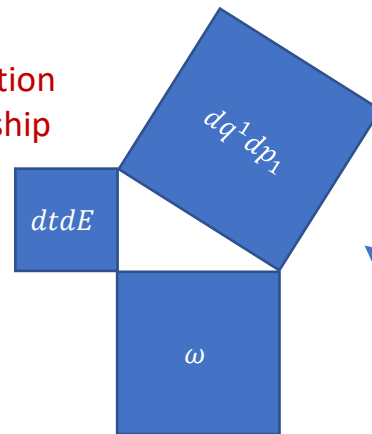
$$\omega = dq^1 dp_1 - dt dE$$

$$dq^1 dp_1 = \omega + dt dE$$



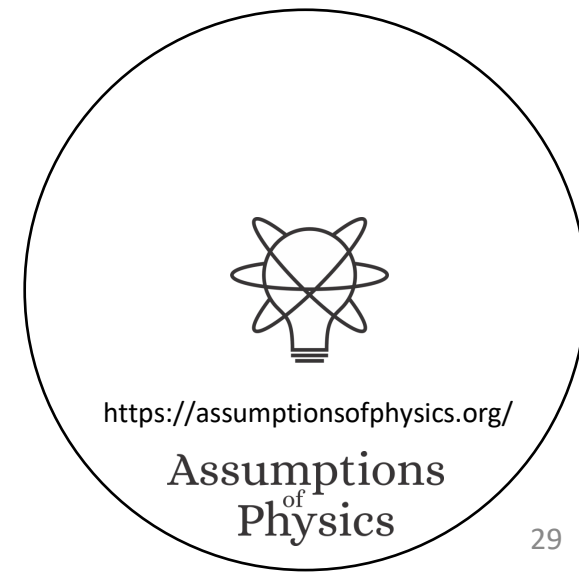
2D representation
of 4D relationship

Independent DOFs are orthogonal



States are counted at
equal time:
temporal DOF
orthogonal to ω

Pythagoras's formula still present



Example: free particle
in an inertial frame

$$\mathcal{H} = \frac{1}{2m} \left(p_i \delta^{ij} p_j - \left(\frac{E}{c} \right)^2 + m^2 c^2 \right)$$

$$\frac{dt}{ds} = \frac{\partial \mathcal{H}}{\partial(-E)} = -\frac{1}{2mc^2} \frac{\partial}{\partial(-E)} (-E)^2 = -\frac{1}{2mc^2} (-2E) = \frac{E}{mc^2}$$

$$\frac{dq^i}{ds} = \frac{\partial \mathcal{H}}{\partial p_i} = \frac{1}{2m} \frac{\partial}{\partial p_i} (p_i \delta^{ij} p_j) = \frac{2\delta^{ij} p_j}{2m} = \frac{p_i}{m}$$

$$\frac{dE}{ds} = -\frac{d(-E)}{ds} = \frac{\partial \mathcal{H}}{\partial t} = 0$$

$$\frac{dp_i}{ds} = -\frac{\partial \mathcal{H}}{\partial q^i} = 0$$

This Hamiltonian constraint is
the generator for proper time
(i.e. affine parameter s
corresponds to proper time)

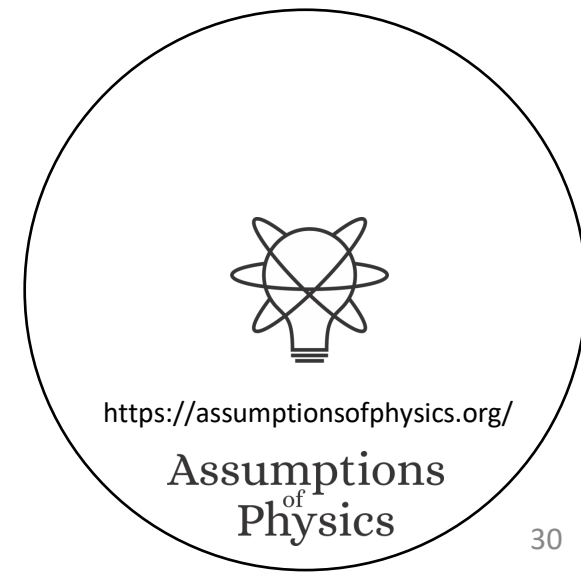
four
momentum

$$\frac{E}{c} = m \frac{d(ct)}{ds}$$

$$p_i = m \frac{dq^i}{ds}$$

four
velocity

proper time!



Compare to “standard” way

Generator for proper time

Nice and quadratic

$$\mathcal{H} = \frac{1}{2m} \left(p_i \delta^{ij} p_j - \left(\frac{E}{c} \right)^2 + m^2 c^2 \right)$$

$$H = c \sqrt{m^2 c^2 + p_i \delta^{ij} p_j}$$

Hamiltonian is time generator

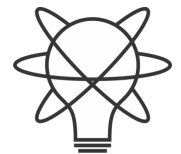
$$\begin{aligned} \frac{dq^i}{dt} &= \frac{\partial H}{\partial p_i} = c \frac{\partial}{\partial p_i} \sqrt{m^2 c^2 + p_i \delta^{ij} p_j} \\ &= \frac{c}{2} \frac{1}{\sqrt{m^2 c^2 + p_i \delta^{ij} p_j}} \frac{\partial}{\partial p_i} (m^2 c^2 + p_i \delta^{ij} p_j) \\ &= \frac{1}{2} \frac{2 \delta^{ij} p_j}{\sqrt{m^2 + \frac{p_i \delta^{ij} p_j}{c^2}}} = \frac{p_i}{\sqrt{m^2 + \frac{|p_i|^2}{c^2}}} \end{aligned}$$

$$\frac{dp_i}{dt} = - \frac{\partial H}{\partial q^i} = 0$$

$$\frac{dq^i}{dt} = \frac{ds}{dt} \frac{dq^i}{ds} = \frac{mc^2}{E} \frac{p_i}{m}$$

Hamiltonian mechanics on the extended phase-space comes out more naturally and is a nicer formulation

Wait... quadratic?
Can E be negative?



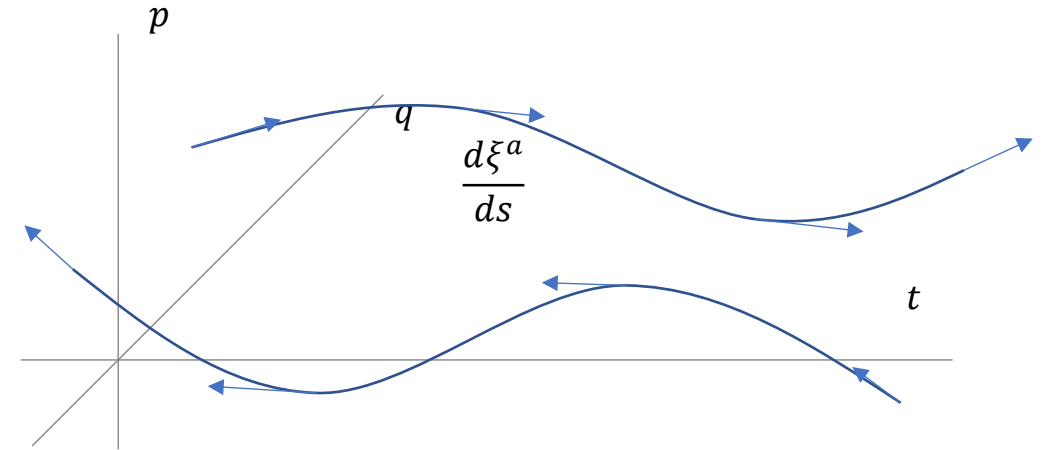
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Assumptions
of
Physics

Classical anti-particles

This can be negative

$$\begin{aligned} d_s t &= \partial_{-E} \mathcal{H} & d_s (-E) &= \partial_t \mathcal{H} \\ d_s q^i &= \partial_{p_i} \mathcal{H} & d_s p_i &= \partial_{q^i} \mathcal{H} \end{aligned}$$



Particle: affine parameter aligned with time

Anti-particle: affine parameter anti-aligned with time

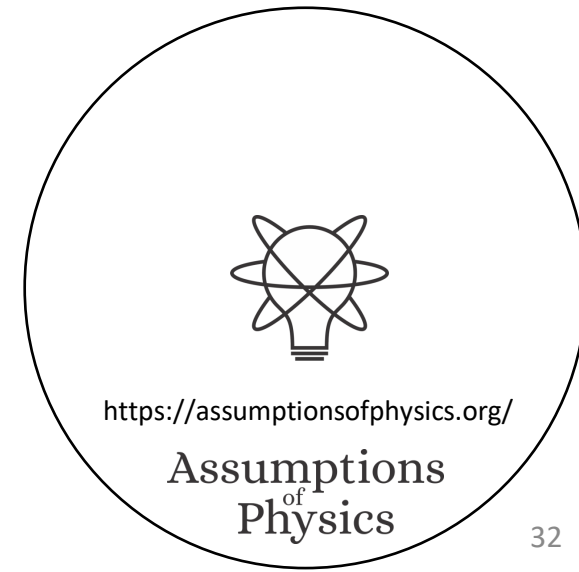
Free particle

$$d_s t = \frac{E}{mc^2}$$

Anti-particle
= negative energy

Parametrization flows
backwards with
respect to time

An evolution cannot
change time alignment



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Assumptions
of
Physics

Recover standard Hamiltonian

Over valid states $\mathcal{H} = 0$ and $E = H(q^i, p_i, t)$ $\mathcal{H} = \lambda(H - E)$

$$E = H$$

$$d_t t = 1 = d_t s d_s t = d_t s \partial_{-E} \mathcal{H} = d_t s \lambda$$

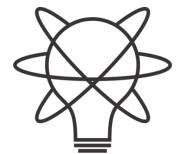
$$d_t s = \frac{1}{\lambda}$$

$$d_t q^i = d_t s d_s q^i = d_t s \partial_{p_i} \mathcal{H} = \frac{1}{\lambda} (\partial_{p_i} \lambda(H - E) + \lambda \partial_{p_i} H) = \partial_{p_i} H$$

$$d_t p_i = d_t s d_s p_i = -d_t s \partial_{q^i} \mathcal{H} = -\frac{1}{\lambda} (\partial_{q^i} \lambda(H - E) + \lambda \partial_{q^i} H) = -\partial_{q^i} H$$

$$d_t E = d_t s d_s E = d_t s \partial_t \mathcal{H} = \frac{1}{\lambda} (\partial_t \lambda(H - E) + \lambda \partial_t H) = \partial_t H$$

Assume no E dependence



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Assumptions
of
Physics

Hamiltonian constraint can “hide” multiple H s

Free particle Hamiltonian constraint

$$\mathcal{H} = \frac{1}{2mc^2} (\sqrt{c^2|p_i|^2 + (mc^2)^2} + E)(\sqrt{c^2|p_i|^2 + (mc^2)^2} - E)$$

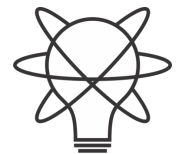
$$\mathcal{H} = (H_1 - E)(H_2 - E) \frac{-1}{2mc^2}$$

$H_1 = -H_2$

Negative energy

Positive energy

$$d_{st} = \frac{1}{2mc^2} (E + E) = \frac{E}{mc^2}$$



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Assumptions
of
Physics

Classical-quantum parallel

Free particle Hamiltonian constraint

$$\mathcal{H} = \frac{1}{2m} (p_\alpha \eta^{\alpha\beta} p_\beta + m^2 c^2) = \frac{1}{2m} (p_i \delta^{ij} p_j - (E/c)^2 + m^2 c^2)$$

Hamiltonian constraint becomes Klein-Gordon equation in QM

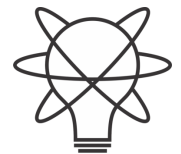
$$\left(\frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \nabla^2 + \frac{m^2 c^2}{\hbar^2} \right) \psi(t, \mathbf{x}) = 0$$

minus sign from i^2

Also note:

$$\mathcal{H}\rho = 0$$

Density can be non-zero
only where $\mathcal{H} = 0$



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Assumptions
of
Physics

Hamiltonian constraint for charged particles

pure four-velocity kinetic term

$$\mathcal{H} = \frac{1}{2m} (p_\alpha - qA_\alpha) g^{\alpha\beta} (p_\beta - qA_\beta) + \frac{1}{2} mc^2 = \frac{1}{2} m u^\alpha g_{\alpha\beta} u^\beta + \frac{1}{2} mc^2$$

$$m u_\alpha = m u^\beta g_{\beta\alpha} = p_\alpha - qA_\alpha$$

Kinetic momentum

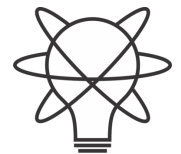
Momentum gauge and EM gauge must cancel out

Compare to EM Klein-Gordon equation

$$[c^2(\partial_\alpha - iqA_\alpha)\eta^{\alpha\beta}(\partial_\beta - iqA_\beta) + m^2c^4]\psi = [c^2D_\alpha\eta^{\alpha\beta}D_\beta + m^2c^4]\psi = 0$$

$$D_\alpha = \partial_\alpha - iqA_\alpha$$

Gauge-covariant derivative is kinetic momentum

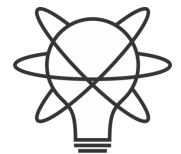


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Assumptions
of
Physics

Takeaways

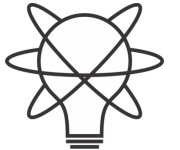
- Relativistic mechanics is recovered without additional assumptions
 - Relativity is needed to make densities/entropy invariant over time transformations
- Hamiltonian mechanics on the extended phase space is a more relativistic formulation
 - Hamiltonian constraint is the generator of the affine parameter, and can be taken to generate proper time for particles and minus proper time for anti-particles
- Multiple connections to quantum mechanics
 - Anti-particles already present in classical Hamiltonian particle mechanics
 - Hamiltonian constraint related to Klein-Gordon (and Dirac) operators
 - Gauge considerations are already present in classical particle mechanics



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Assumptions
of
Physics

Directional DOF

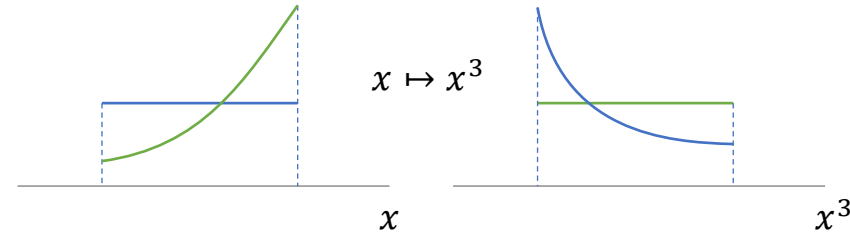


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Assumptions
of
Physics

Why DOFs are charted by two conjugate variables

Density, entropy, uniform distributions
NOT in general coordinate-invariant



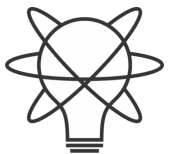
The principle of indifference fails over continuous quantities

Phase space (symplectic) structure is the
only one that supports coordinate-
invariant density, entropy, state count

Each unit variable (i.e. coordinate) paired
with a conjugate of inverse units: number
of states $\Delta q \Delta k$ is invariant

$\Delta k = 1 \text{ m}^{-1}$	$\hat{q} = 100 \text{ cm/m } q$		$\Delta \hat{k} = 0.01 \text{ cm}^{-1}$
	1	1	
	$\Delta q = 1 \text{ m}$	$\Delta \hat{q} = 100 \text{ cm}$	

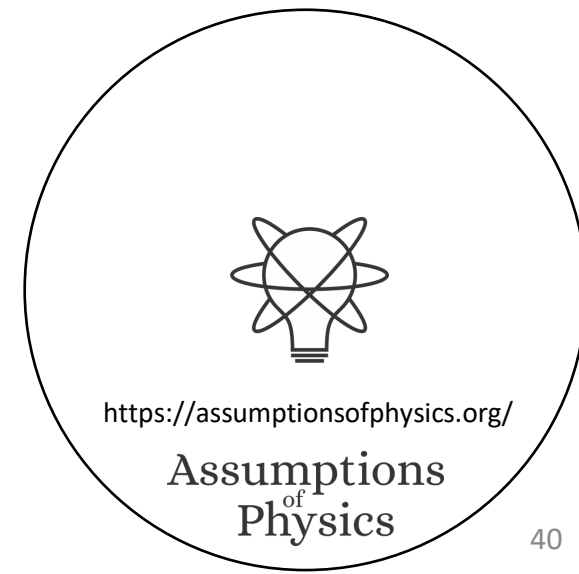
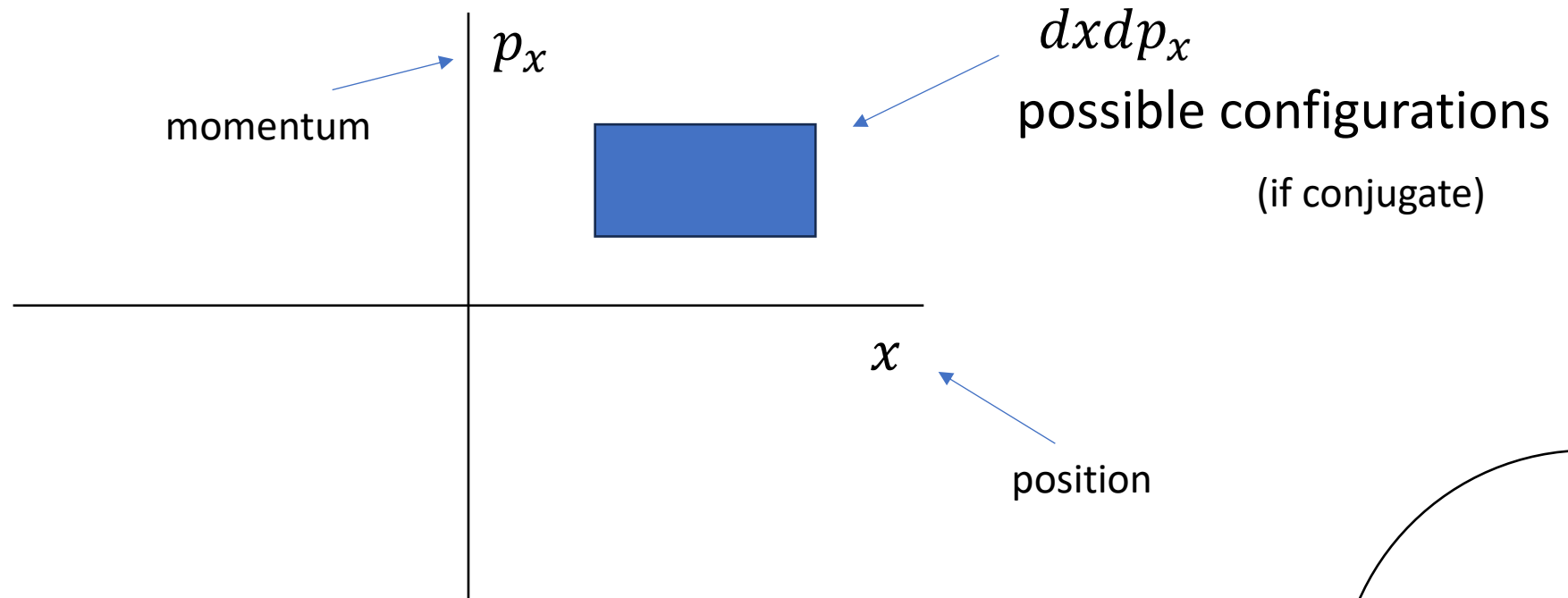
Conjugate quantities recover frame independence of states



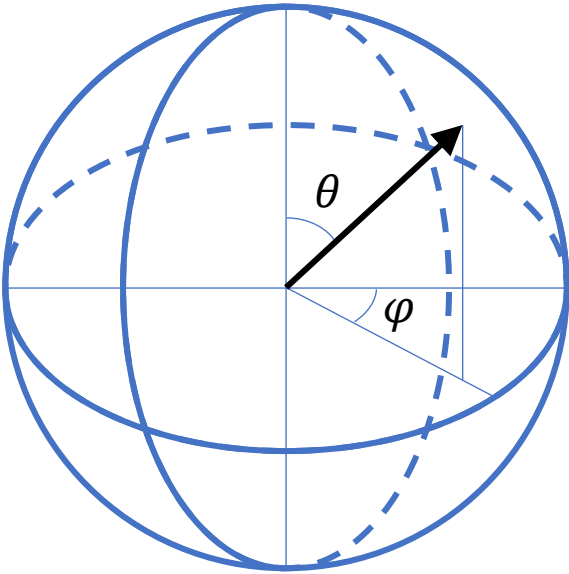
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Assumptions
of
Physics

Each DOF is charted by two variables

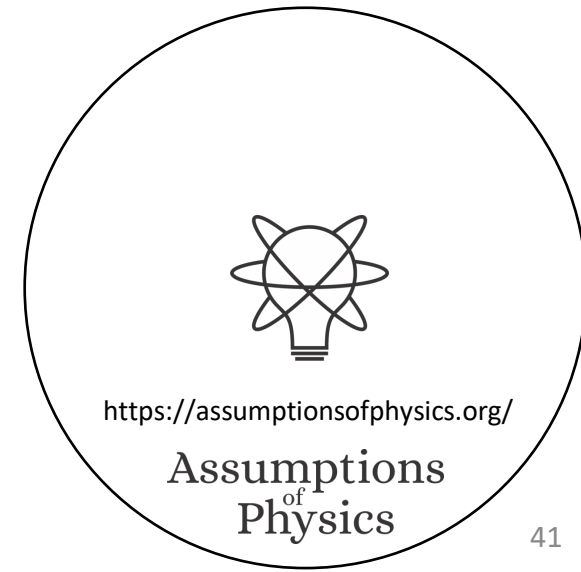


Angular momentum

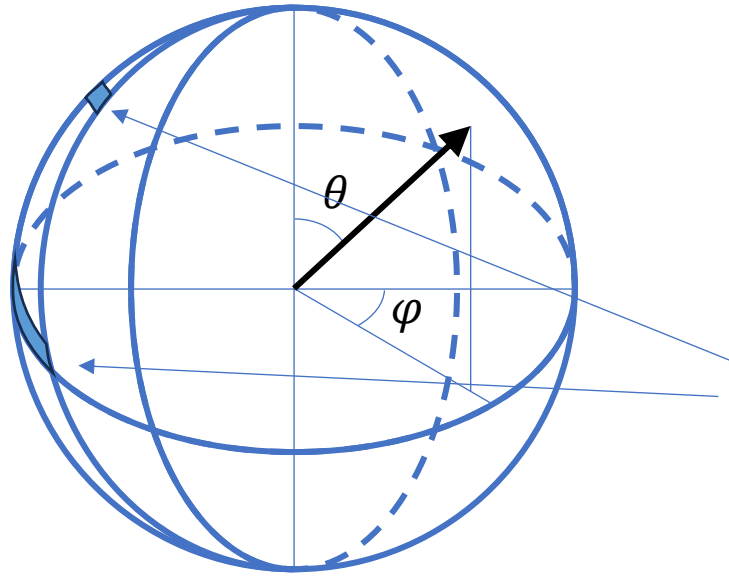


$$\begin{aligned}\vec{L} &= [L_x, L_y, L_z] \\ &= [L \sin \theta \cos \varphi, L \sin \theta \sin \varphi, L \cos \theta]\end{aligned}$$

Direction independent
from the magnitude



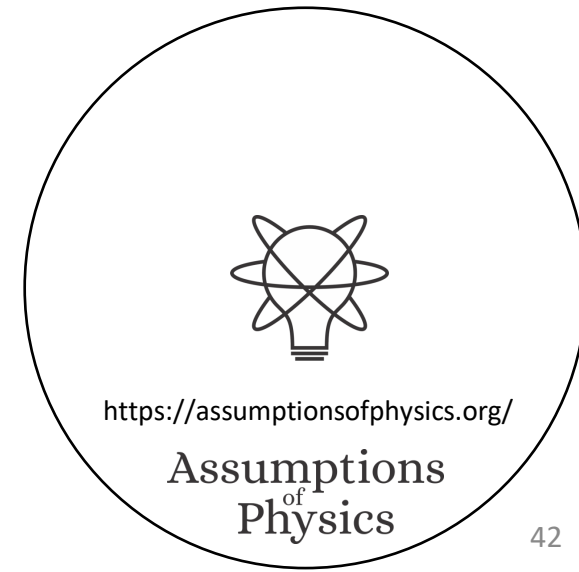
Angular momentum direction



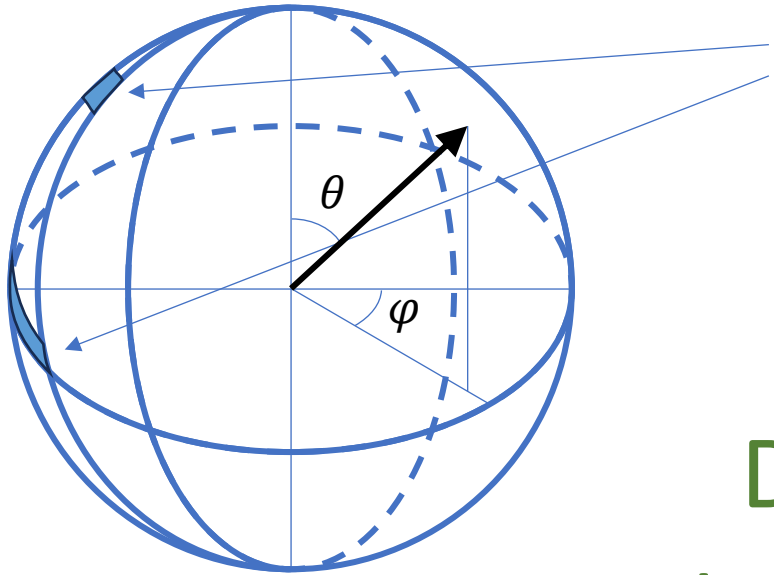
Direction is an independent DOF

However, θ and φ are not conjugate

$d\theta d\varphi \neq$ possible configurations



Angular momentum direction

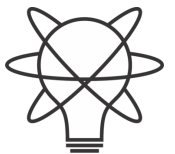


$$Ld\Omega = L \sin \theta d\theta d\varphi = d(L \cos \theta) d\varphi \\ = dL_z d\varphi_{xy}$$

φ_{xy} and L_z are conjugate

Direction is charted by
two conjugate quantities

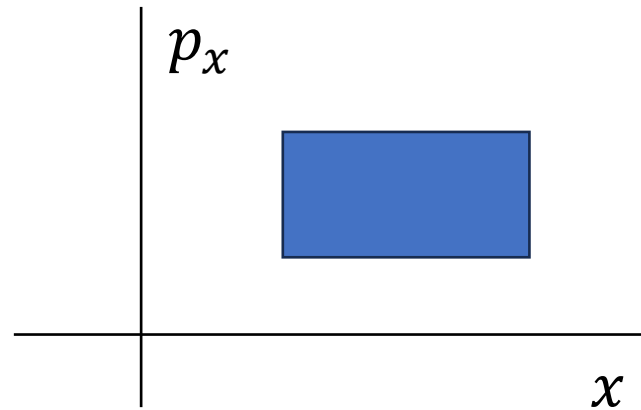
$\Leftrightarrow$ Space is three dimensional



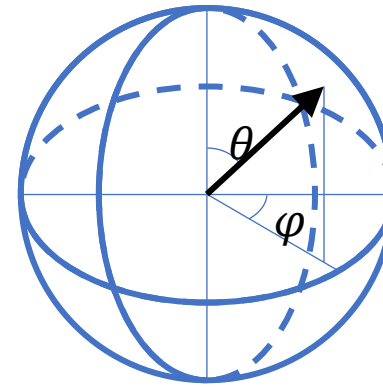
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Assumptions
of
Physics

Independent DOF

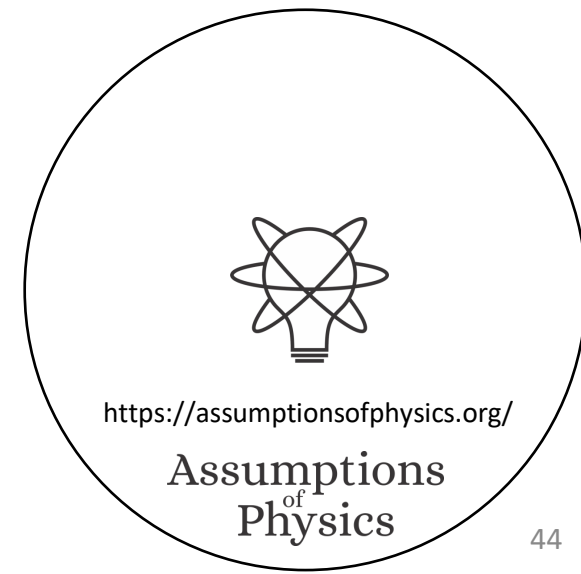


Two dimensional



Independent
directional DOF

Space three
dimensional

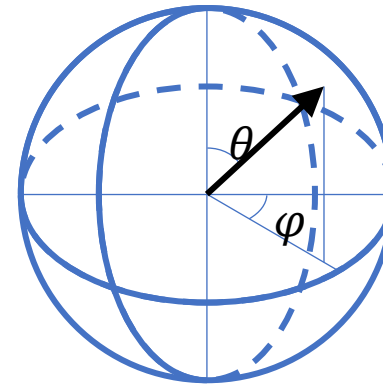
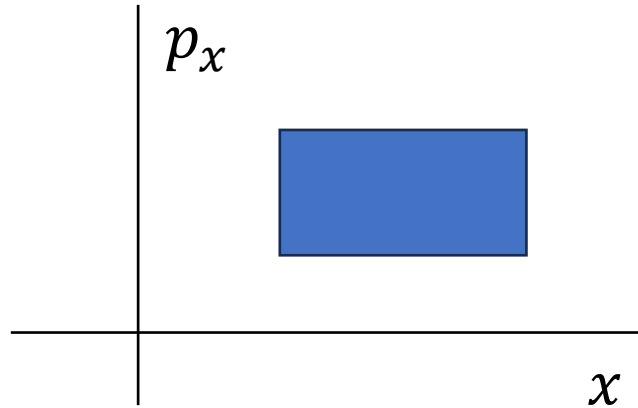


Independent DOF

Suppose



Two dimensional



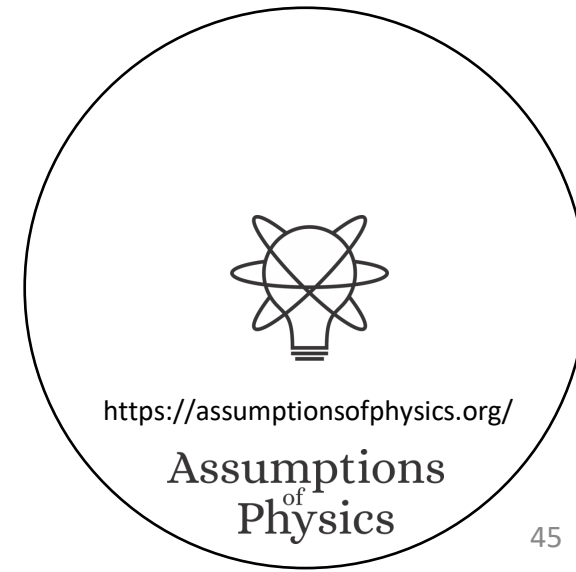
Independent directional DOF

???



Space three dimensional

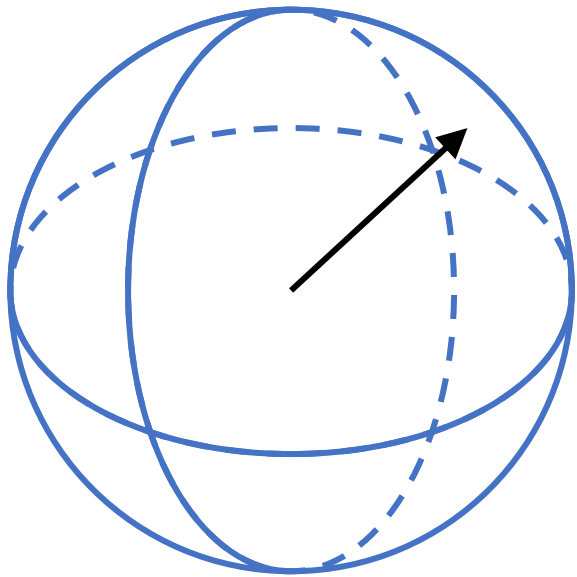
A single directional DOF could be made
of multiple independent DOF



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Assumptions
of
Physics

Direction in multiple dimensions



Direction is a point on a hypersphere

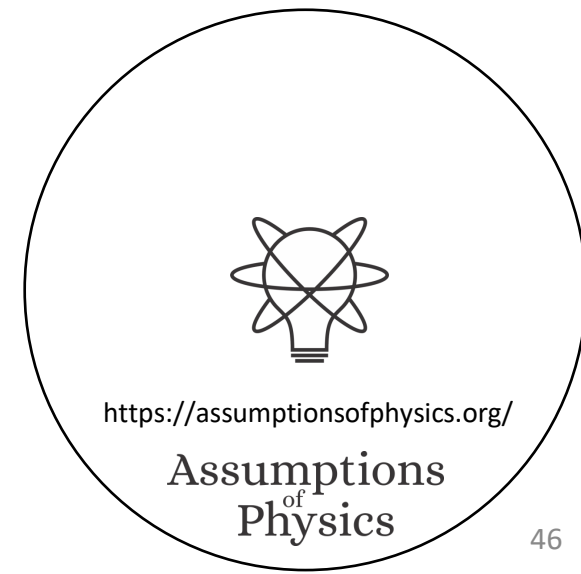
$$\hat{L} = [\hat{L}_1, \hat{L}_2, \hat{L}_3, \dots, \hat{L}_n]$$

$$\sum_i \hat{L}_i = 1$$

normalized vector

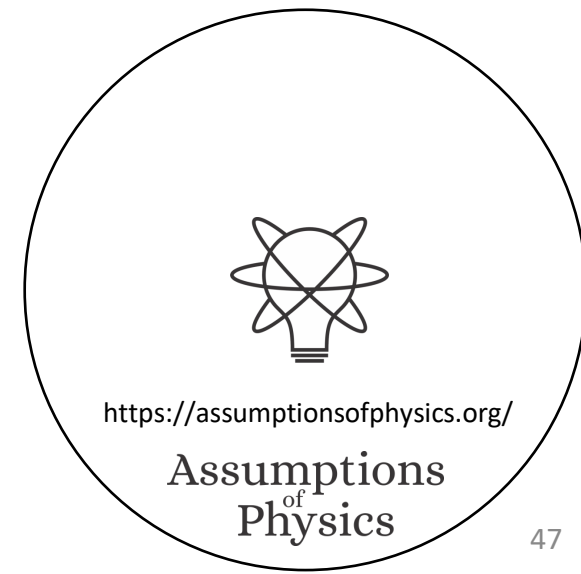
2-sphere is the only symplectic manifold

Direction cannot be described
by multiple independent DOFs



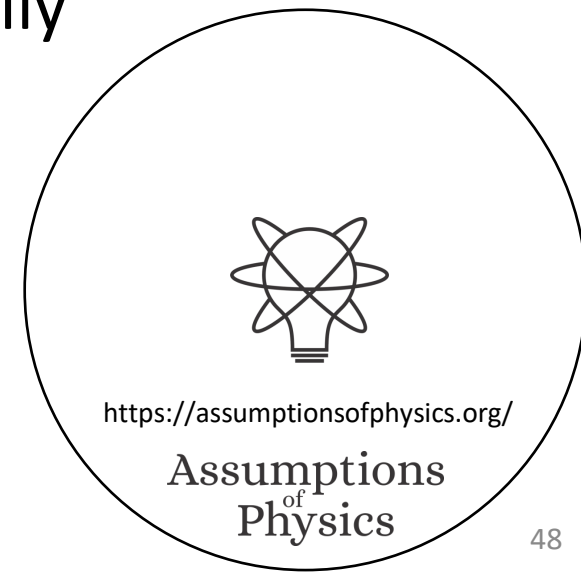
Takeaways

- Invariant densities, entropy and count of states cannot be defined in general over the continuum
- The additional structure required is exactly the structure of phase space (symplectic structure)
- If a direction in space forms an independent DOF, then it must be a sphere that has the structure of phase space
- But the two-sphere is the only one possible
- $\Rightarrow$ Space has three dimensions



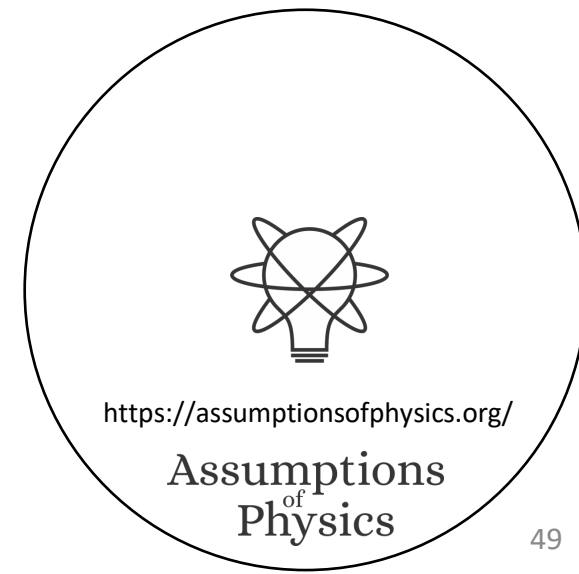
Wrapping it up

- All the core elements of classical mechanics can be rederived and understood using four physically meaningful assumptions
 - Infinitesimal Reducibility (IR), DOF independence (IND), Determinism and Reversibility (DR) and Kinematic Equivalence (KE)
- Many physical ideas are a consequence of those assumptions
 - Principle of stationary action, massive particles under scalar/vector potential forces, relativistic mechanics, phase space (conjugate quantities), ...
- A great number of better insights from a theory that is generally considered as “understood”
 - Probably even better insights to be found



Reading group?

- If there is interest, we can set up an online reading group, livestreamed on the research channel
 - Go through the book chapter on classical mechanics together, section by section, discuss and think of potential diagrams/figures that can help
- I'd need one person to volunteer to organize
 - Help get list of interested people, find a time, send reminders, prepare thumbnail, etc...



To learn more

- Project website

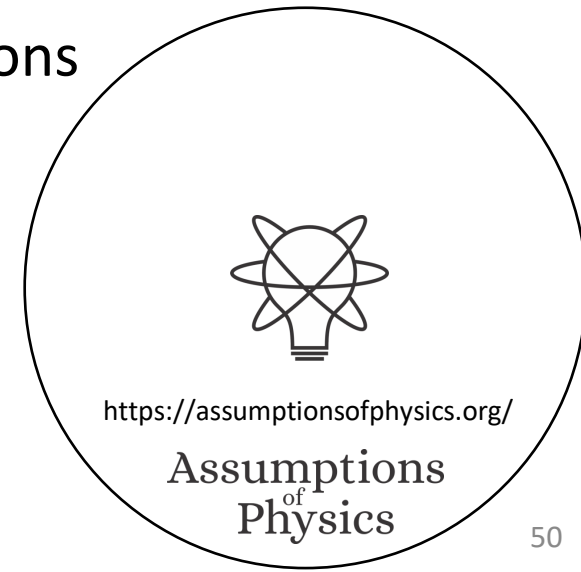
- <https://assumptionsofphysics.org> for papers, presentations, ...
- <https://assumptionsofphysics.org/book> for our open access book (updated every few years with new results)

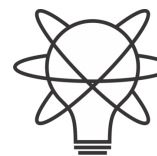
- YouTube channels

- <https://www.youtube.com/@gcarcassi>
Videos with results and insights from the research
- <https://www.youtube.com/@AssumptionsofPhysicsResearch>
Research channel, with open questions and livestreamed work sessions

- GitHub

- <https://github.com/assumptionsofphysics>
Book, research papers, slides for videos...





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